


22, 23 y 24 de Octubre de 2008. El Escorial (Madrid)

Cuarta reunión de la red temática Dance

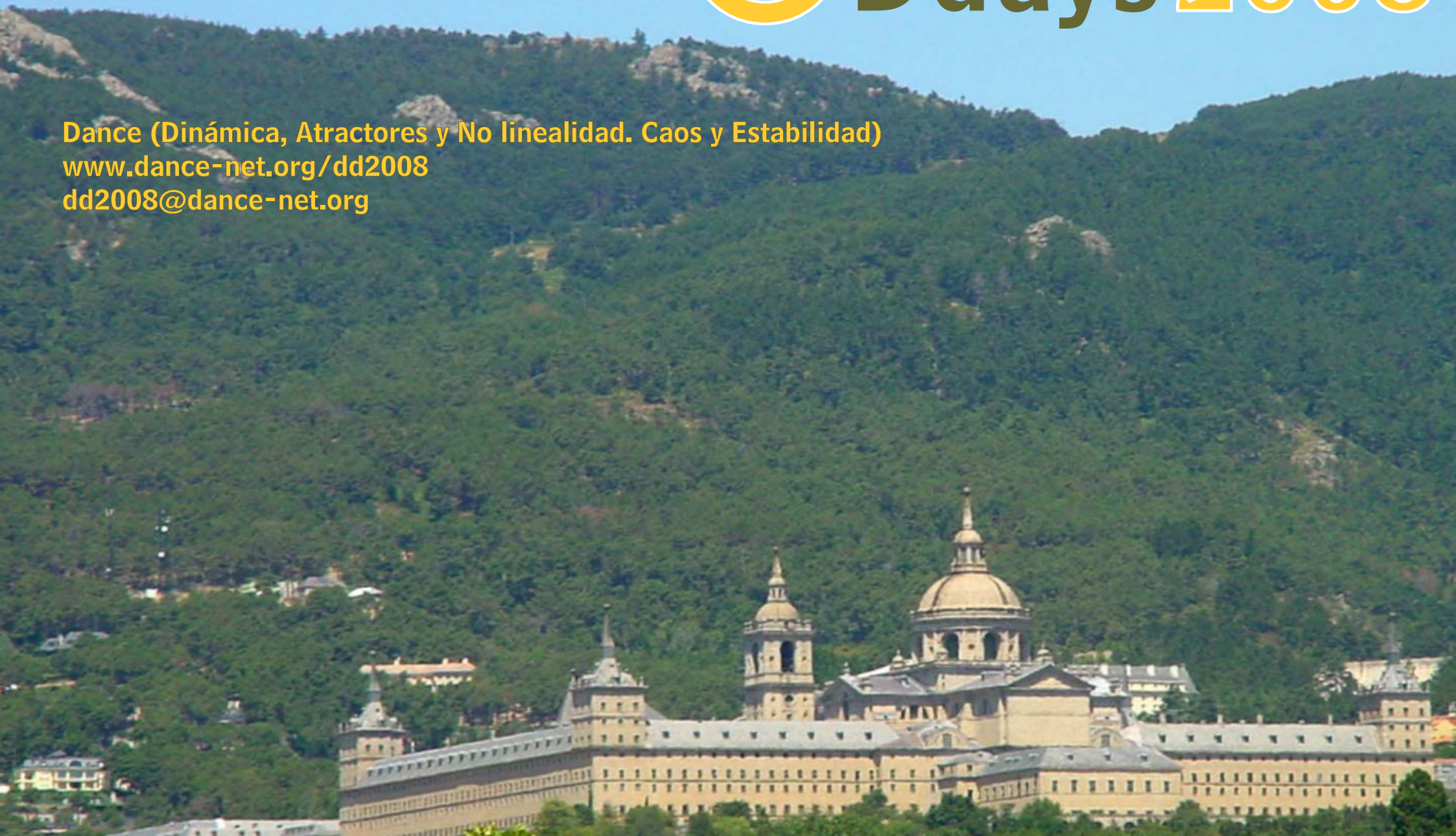
The logo for Ddays 2008 features a stylized lowercase 'd' in yellow with a white outline. Inside the 'd' is a yellow circle with a white outline, resembling a target or a stylized 'o'. To the right of this symbol, the word 'Ddays' is written in a bold, dark grey sans-serif font, and '2008' is written in a yellow sans-serif font with a white outline.

# Ddays 2008

Dance (Dinámica, Atractores y No linealidad. Caos y Estabilidad)

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# Multistep cosine methods for second-order partial differential equations

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**AIM:** Efficient numerical integration of PDEs of second-order in time with qualitative properties to imitate

$$u_{tt}(x, t) = Au(x, t) + g(t, u(x, t)), \quad x \in [a, b], \quad 0 < t < T$$

$$u(x, 0) = u_0(x)$$

$$u_t(x, 0) = v_0(x)$$

periodic boundary conditions

$Au = u_{xx} \rightarrow$  nonlinear wave equation

$Au = -u_{xxxx} \rightarrow$  Euler-Bernoulli equation

**Important assumption:** Enough regularity of the solution

## Literature

- Integration of highly oscillatory problems
  - energy of the solution being bounded
  - stiff part easy to integrate

### Gautschi-type methods

Hochbruch&Lubich, Numer. Math (1999)

Hairer & Lubich, SIAM J. Numer. (2000)

Grimm & Hochbruch, J. Phys. A Math. Gen. (2006)

### Mollified impulse methods

García-Archilla, Sanz-Serna, Skeel, SIAM J. Sci. Comp. (1998)

Sanz-Serna, SIAM J. Numer. (2008)

second-order accurate

## Literature

- Integration of parabolic PDEs with exponential integrators

stiff and linear part integrated exactly

→ explicit and stable methods of arbitrary high order

Norsett, Lecture Notes (1969)

Cox & Matthews, JCP (2002)

Hochbruch & Ostermann, APNUM (2005)

Hochbruch & Ostermann, SIAM J. Num. Anal. (2005)

Calvo & Palencia, Numer. Math. (2006)

no qualitative properties to imitate

We will use symmetric explicit multistep cosine methods

explicit multistep  $\implies$  1 funct.eval./step

cosine  $\implies$  stiff part integrated exactly

symmetric  $\implies$  qualitative properties are conserved



Literature on symmetric multistep methods for ODEs of second-order in time  $\ddot{y} = f(y)$

$$\rho(E)y_n = h^2\sigma(E)f(y_n), \quad \rho(x) = x^k\rho\left(\frac{1}{x}\right), \quad \sigma(x) = x^k\sigma\left(\frac{1}{x}\right).$$

If 1 is the only double root of  $\rho$   $\longrightarrow$  good qualitative behaviour

In another case  $\longrightarrow$  exponential error growth with time in general

Cano & Sanz-Serna, IMA J. Numer. Anal. (1998)

Hairer & Lubich, Numer. Math. (2004)

Cano, Numer. Math. (2006)  $\longrightarrow$  generalization to PDEs

After pseudospectral discretization in space,

$$\ddot{Y}(t) = -\Omega^2 Y(t) + G(t, Y(t)), \quad \Omega \rightarrow \text{diagonal and real}$$

### Construction of methods

$$Y(t_n + h) = \cos(h\Omega)Y(t_n) + h\text{sinc}(h\Omega)\dot{Y}(t_n) + \int_0^h \Omega^{-1} \sin((h-s)\Omega)G(t_n + s, Y(t_n + s))ds$$

$h \rightarrow -h$  and adding

$$\begin{aligned} Y(t_n + h) - 2\cos(h\Omega)Y(t_n) + Y(t_n - h) \\ = \int_0^h \Omega^{-1} \sin((h-s)\Omega)[G(t_n + s, Y(t_n + s)) + G(t_n - s, Y(t_n - s))]ds. \end{aligned}$$

When  $G(t_n + s, Y(t_n + s)) \approx G(t_n, Y_n) \rightarrow$  Gautschi method

$$Y_{n+2} - 2\cos(h\Omega)Y_{n+1} + Y_n = h^2\gamma_1(h\Omega)G_{n+1}, \quad \gamma_1(\epsilon) = 2\epsilon^{-2}(1 - \cos(\epsilon)),$$



If  $G(t_n+s, Y(t_n+s)) \approx \sum_{l=0}^{k-1} G(t_{n-l}, Y_{n-l}) L_{-l,h}(s) \rightarrow$  higher order method

$L_{-l,h}$ : Lagrange polynomial  $\{0, -h, \dots, -lh, \dots, -(k-1)h\}$

$k = 3$       MC3

$$Y_{n+3} - 2 \cos(h\Omega) Y_{n+2} + Y_{n+1} = h^2 [\gamma_2(h\Omega) G_{n+2} + \gamma_1(h\Omega) G_{n+1} + \gamma_0(h\Omega) G_n]$$

$$\gamma_0(\epsilon) = \epsilon^{-2} [1 - 2\epsilon^{-2} + 2\epsilon^{-2} \cos(\epsilon)].$$

$$\gamma_1(\epsilon) = \epsilon^{-2} [-2 + 4\epsilon^{-2} - 4\epsilon^{-2} \cos(\epsilon)],$$

$$\gamma_2(\epsilon) = \epsilon^{-2} [3 - 2 \cos(\epsilon) - 2\epsilon^{-2} + 2\epsilon^{-2} \cos(\epsilon)].$$

$\Omega \rightarrow 0$ : 3rd-order explicit Störmer multistep method, not symmetric

## SMC4

$$\begin{aligned}
 & Y(t_n + 2h) - \delta(h\Omega)Y(t_n + h) - 2(\cos(2h\Omega) - \delta(h\Omega)\cos(h\Omega))Y(t_n) - \delta(h\Omega)Y(t_n - h) + Y(t_n - 2h) \\
 &= \int_0^{2h} \Omega^{-1} \sin((2h-s)\Omega)[G(t_n + s, Y(t_n + s)) + G(t_n - s, Y(t_n - s))]ds \\
 &\quad - \delta(h\Omega) \int_0^h \Omega^{-1} \sin((h-s)\Omega)[G(t_n + s, Y(t_n + s)) + G(t_n - s, Y(t_n - s))]ds.
 \end{aligned}$$

Interpolation in  $\{-h, 0, h\} \rightarrow$  symmetry

$$\delta(\epsilon) = \cos(\epsilon) \begin{cases} \rho_0(x) = x^4 - x^3 - x + 1 = (x-1)^2(x^2 + x + 1) \\ \text{stability} \rightarrow \text{roots of } \rho_\epsilon(x) \text{ of unit modulus} \end{cases}$$

$$\begin{aligned}
 x_1(\epsilon) &= \frac{-\cos(\epsilon) + i\sqrt{4 - \cos^2(\epsilon)}}{2}, & x_2(\epsilon) &= \frac{-\cos(\epsilon) - i\sqrt{4 - \cos^2(\epsilon)}}{2}, \\
 x_3(\epsilon) &= \cos(\epsilon) + i\sin(\epsilon), & x_4(\epsilon) &= \cos(\epsilon) - i\sin(\epsilon).
 \end{aligned}$$

$$\begin{aligned}
 & Y_{n+4} - \cos(h\Omega)Y_{n+3} + (-2\cos(2h\Omega) + 2\cos(h\Omega)^2)Y_{n+2} - \cos(h\Omega)Y_{n+1} + Y_n \\
 &= h^2[\gamma_1(h\Omega)G_{n+3} + \gamma_2(h\Omega)G_{n+2} + \gamma_1(h\Omega)G_{n+1}],
 \end{aligned}$$

$$\begin{aligned}
 \gamma_1(\epsilon) &= \epsilon^{-2}(4 - \cos(\epsilon)) + 2\epsilon^{-4}[\cos(2\epsilon) - 1 + \cos(\epsilon)(1 - \cos(\epsilon))] \\
 \gamma_2(\epsilon) &= 2\epsilon^{-2}(-3 - \cos(2\epsilon) + \cos^2(\epsilon)) + 4\epsilon^{-4}(1 - \cos(2\epsilon) - \cos(\epsilon)(1 - \cos(\epsilon)))
 \end{aligned}$$

## Local truncation error ( $d_n$ )

Gautschi  $\rightarrow q = 2$

MC3  $\rightarrow q = 3$

SMC4  $\rightarrow q = 4$

By using error of interpolation

- If  $\frac{d^q}{dt^q}G(t, Y(t))$  exists and  $\left| \left( \frac{d^q}{dt^q}G(t, Y(t)) \right)_\lambda \right| \leq K_{\lambda, q}$

$$|d_{n, \lambda}| \leq C_q K_{\lambda, q} h^{q+2}, \quad C_2 = 1, \quad C_3 = (1 + 2\sqrt{3}/5)/2, \quad C_4 = (4 + \frac{\sqrt{3}}{54})$$

- Under any of the following assumptions of regularity

(i)  $\frac{d^q}{dt^q}g(t, u(t, x))$  can be extended to a holomorphic function on a complex band  $|\text{Im}(z)| \leq \hat{B}$  for  $0 \leq t \leq T$ , and there  $|\frac{d^q}{dt^q}g(t, u(t, z))| \leq \hat{C}$ ,

(ii)  $\frac{d^q}{dt^q}g(t, u(t, x))$  admits  $m \geq 0$  continuous and periodic derivatives with respect to  $x$  and  $\frac{d^{m+1}}{dx^{m+1}} \frac{d^q}{dt^q}g(t, u(t, x))$  exists and is piecewise  $C^1$  in  $[a, b]$ .

$$\|d_n\| = O(h^{q+2})$$

## Global error

$$\bar{Y}_n = [Y_{n+k-1}, \dots, Y_n]^T, \quad \bar{G}(\bar{Y}_n) = [G(Y_{n+k-1}), \dots, G(Y_n)]^T$$

$$\bar{Y}_{n+1} = R(h\Omega)\bar{Y}_n + h^2 B(h\Omega)\bar{G}(\bar{Y}_n), \quad B(\epsilon) = \begin{pmatrix} \gamma_{k-1}(\epsilon) & \dots & \gamma_1(\epsilon) & 0 \\ 0 & 0 & \dots & 0 \\ \vdots & & \ddots & \vdots \\ 0 & \dots & \dots & 0 \end{pmatrix},$$

$$\text{Gautschi} \rightarrow R(\epsilon) = \begin{pmatrix} 2 \cos(\epsilon) & -1 \\ 1 & 0 \end{pmatrix}, \quad \text{MC3} \rightarrow R(\epsilon) = \begin{pmatrix} 2 \cos(\epsilon) & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\text{SMC4} \rightarrow R(\epsilon) = \begin{pmatrix} \cos(\epsilon) & 2 \cos(2\epsilon) - 2 \cos(\epsilon)^2 & \cos(\epsilon) & -1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

$$\|B(h\Omega)\| \leq \sup_{\lambda \in \sigma(\Omega)} \|B(h\lambda)\|$$

Stability  $\|R^j(\epsilon)\| \leq \bar{C}_j, \quad \epsilon \in \mathbb{R}$

Eigenvalues of  $R(\epsilon)$   $\left\{ \begin{array}{l} \text{of modulus } \leq 1 \\ \text{those of unit modulus} \\ \text{at most double} \end{array} \right\} \rightarrow$  Schur decomposition

Gautschi  $\epsilon = 2m\pi, \pi + 2m\pi, \rightarrow x_{1,2} = e^{\pm i\epsilon} = +1, -1$

	$\epsilon = 2m\pi$	$\epsilon = \frac{\pi}{2} + 2m\pi$	$\epsilon = 3\frac{\pi}{2} + 2m\pi$	$\epsilon = \pi + 2m\pi$
$x_1$	$\frac{-1+i\sqrt{3}}{2}$	$i$	$i$	$\frac{1+i\sqrt{3}}{2}$
$x_2$	$\frac{-1-i\sqrt{3}}{2}$	$-i$	$-i$	$\frac{1-i\sqrt{3}}{2}$
$x_3$	$1$	$i$	$-i$	$-1$
$x_4$	$1$	$-i$	$i$	$-1$

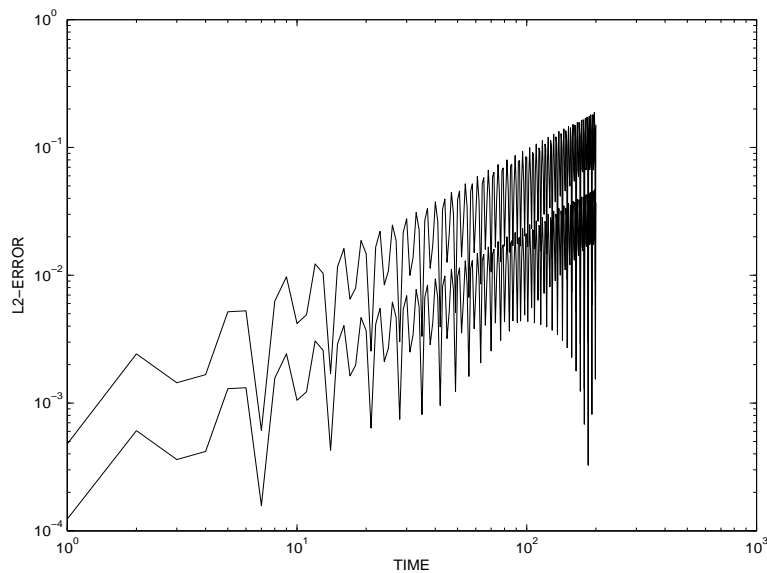
- Under the mentioned assumptions of **regularity**, whenever  $Y(t_\nu) - Y_\nu = O(h^{q+1}), \nu = 0, 1, \dots, k-1, Y(t_n) - Y_n = O(h^q), 0 \leq nh \leq T.$

## Growth of error with time for the nonlinear wave equation

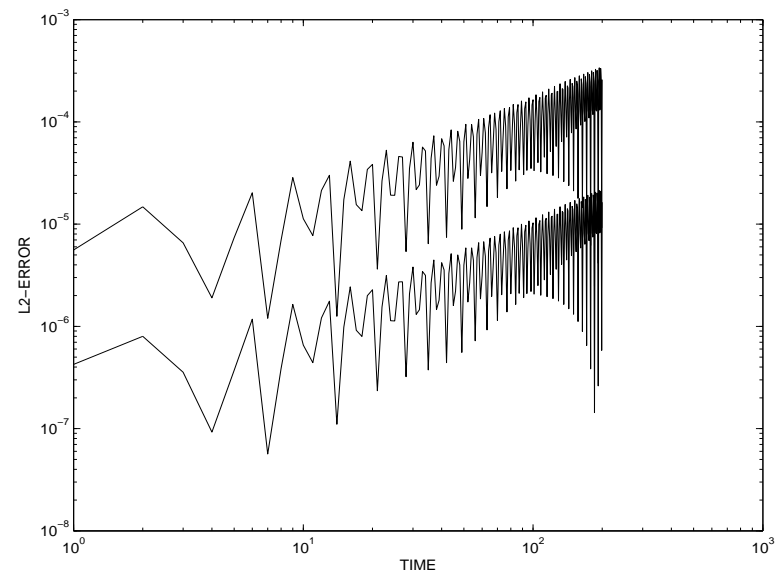
Pseudospectral discretization in space ( $M = 250$ )

$$\Omega = \text{diag}(\lambda_k)_{k=-M}^M = \text{diag}\left(\frac{2\pi k}{b-a}\right)_{k=-M}^M = \text{diag}\left(\frac{2\pi k}{100}\right)_{k=-M}^M$$

Gautschi



SMC4



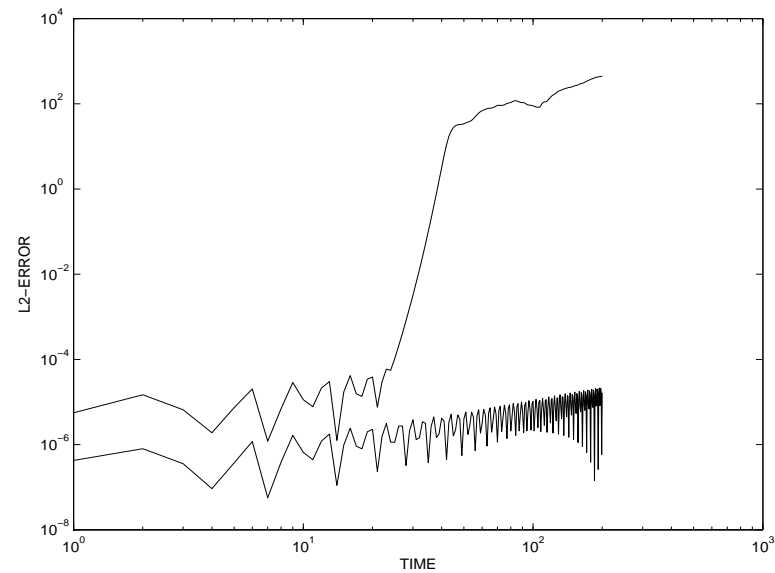
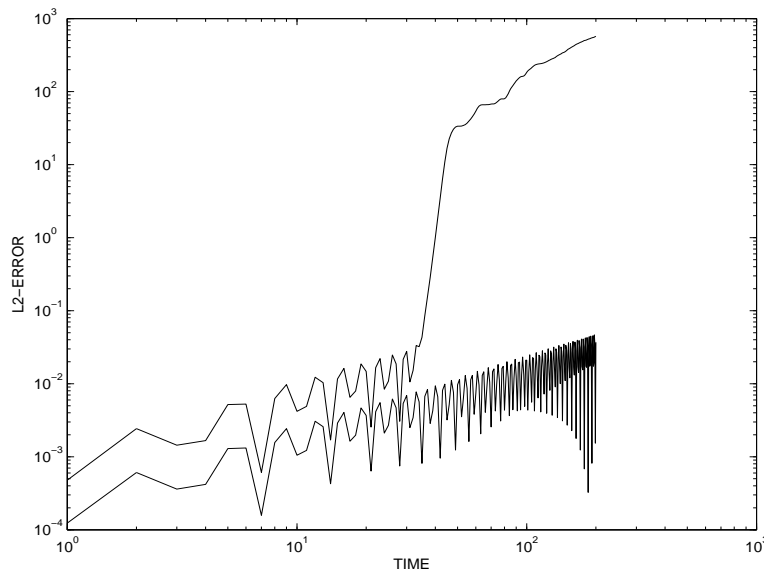
$$h = 1/10, 1/20,$$

$$|\lambda_k h| \leq 5\pi h \leq \pi/2$$

## Resonances

$$h = 1/10, 1/20,$$

$$M = 500$$



$h = 1/10, \lambda_M h = \pi \rightarrow$  leads to double root of  $\rho_\epsilon(x)$

$$h = 1/20, |\lambda_k h| \leq \pi/2$$



## Explanation for resonances

Simplified problem  $\ddot{y}(t) = -\lambda^2 y(t) - y(t)$

Numerical solution for stepsize  $h$ ,  $Y_n = \sum_{j=1}^k \delta_j [x_j(\lambda h, h)]^n$

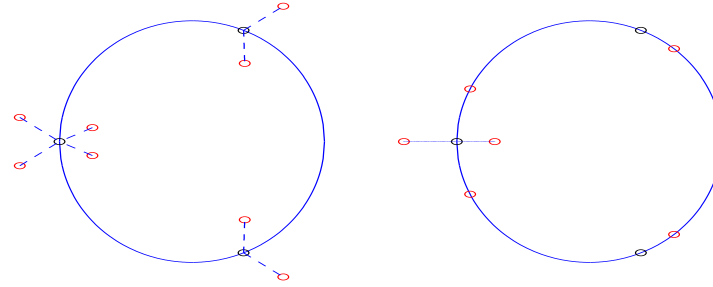
$x_j(\epsilon, h) \rightarrow$  roots of  $r_{\epsilon, h}(x) \equiv \rho_\epsilon(x) + h^2 \sigma_\epsilon(x)$

$$\rho_\epsilon(x) = \alpha_k(\epsilon)x^k + \cdots + \alpha_0(\epsilon), \quad \sigma_\epsilon(x) = \gamma_{k-1}(\epsilon)x^{k-1} + \cdots + \gamma_1(\epsilon),$$

Symmetric methods  $\left\{ \begin{array}{l} \alpha_j(\epsilon) = \alpha_{k-j}(\epsilon), \quad \gamma_j(\epsilon) = \gamma_{k-j}(\epsilon), \\ \rho_\epsilon(x) = x^k \rho_\epsilon(1/x), \quad \sigma_\epsilon(x) = x^k \sigma_\epsilon(1/x), \end{array} \right.$

$$x_j(\epsilon, h) \text{ root of } r_{\epsilon, h}(x) \implies \left\{ \begin{array}{l} \overline{x_j(\epsilon, h)} \\ 1/x_j(\epsilon, h) \end{array} \right\} \text{ also roots of } r_{\epsilon, h}(x)$$

Stable methods  $\rightarrow$  roots of unit modulus of  $r_{\epsilon, 0}(x)$  at most double



Gautschi      small  $h$

• If  $\epsilon$  far from  $m\pi$  ( $m \in \mathbb{Z}$ ),  $|x_j(\epsilon, h)| = 1$ ,  $j = 1, 2$ .

• If  $\epsilon$  near  $m\pi$  ( $m \in \mathbb{Z}$ ),

$|x_j(\epsilon, h)| = 1$ ,  $j = 1, 2$ .

or

$\left. \begin{array}{l} x_1(\epsilon, h) \text{ real and } |x_1(\epsilon, h)| > 1, \\ x_2(\epsilon, h) = 1/x_1(\epsilon, h) \end{array} \right\} \implies \text{resonance}$

By using Taylor series expansion of  $r_{\epsilon, h}(x)$  around  $\epsilon = \epsilon_l$ ,  
 $h = 0$ ,  $x = x_l$  with  $x_l$  double root of  $r_{\epsilon_l, 0}$ ,  $\epsilon = \epsilon_l + \eta$ ,  $x = x_l + r$

$\epsilon_l = m\pi \left\{ \begin{array}{l} \text{even } m \implies r \text{ purely imaginary in first approximation} \\ \text{odd } m \implies r \text{ real} \implies \text{resonance} \end{array} \right.$

## SMC4

small  $h$

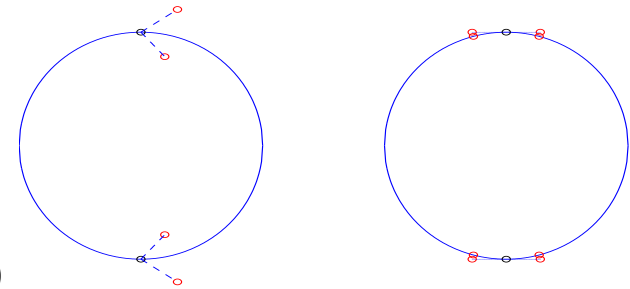
- If  $\epsilon$  far from  $m\pi, \pi/2 + m\pi$  ( $m \in \mathbb{Z}$ ),  $|x_j(\epsilon, h)| = 1$ .
- If  $\epsilon$  near  $m\pi$ ,  $|x_j(\epsilon, h)| = 1$ , or
 
$$\left. \begin{array}{l} x_1(\epsilon, h) \text{ real and } |x_1(\epsilon, h)| > 1, \\ x_2(\epsilon, h) = 1/x_1(\epsilon, h) \end{array} \right\} \implies \text{resonance}$$

By Taylor expansion around  $\epsilon = \epsilon_l, h = 0, x = x_l, \epsilon = \epsilon_l + \eta, x = x_l + r,$

- $$\left\{ \begin{array}{l} \text{even } m \implies r \text{ purely imaginary} \implies \text{no resonance} \\ \text{odd } m \implies r \text{ real} \implies \text{resonance} \end{array} \right.$$

- If  $\epsilon$  near  $\pi/2 + m\pi$ , by Taylor expansion,

$$r = (-1)^m \frac{2\eta \pm \sqrt{4\eta^2 + 16(ch^2 + 2\eta^2)}}{8}, \quad c > 0$$



- real  $\left\{ \begin{array}{l} \text{positive} \\ \text{negative} \end{array} \right. \implies |x_j(\epsilon, h)| = 1 \implies \text{no resonance}$

## Filters to avoid resonances

- Annihilating components  $\rightarrow$  resonant frequencies

$$\ddot{Y}(t) = -\Omega^2 Y(t) + G(t, Y(t)), \quad \Omega = \text{diag}(\lambda_r),$$

If  $|\lambda_r h - \epsilon_l| < \eta_0(\epsilon_l)$  then  $Y_r = 0$

$\epsilon_l$  one of the values which leads to resonance  
 $\eta_0(\epsilon_l)$  deduced from Taylor expansion before

Under our assumptions of regularity, convergence conserved

- Standard filters: filter function  $\phi$

$$\sum_{l=0}^k \alpha_l(h\Omega) Y_{n+l} = h^2 \sum_{l=0}^k \gamma_l(h\Omega) G(\phi(h\Omega) Y_{n+l}), \quad n \geq 0$$

When applied to linear problem  $\ddot{y}(t) = -\lambda^2 y(t) - y(t)$ ,  
control of roots of  $\tilde{r}_{\epsilon, h}(x) = \rho_{\epsilon}(x) + h^2 \tilde{\sigma}_{\epsilon}(x)$ ,  $\tilde{\sigma}_{\epsilon}(x) = \phi(\epsilon) \sigma_{\epsilon}(x)$ .

## Gautschi method

- To avoid resonance in simplified problem,

$$\phi(0) \geq 0,$$

$$\phi(\pi + 2m\pi) \leq 0, \quad m \in \mathbb{Z}, \quad (\text{Taylor expansion})$$

$$\text{If } \phi(\pi + 2m\pi) = 0 \text{ then } \phi'(\pi + 2m\pi) = 0.$$

- To conserve consistency of order 2 for the more general problem, with the assumptions of regularity,  $g$  Lipschitz in  $L^2$ -norm on its second variable and  $\|Au\|_2$  bounded,

$$\frac{\phi(\epsilon) - 1}{\epsilon^2} \text{ bounded for real } \epsilon \quad (\text{local truncation error revisited})$$

$$\text{Possible filter functions: } \left(\frac{\sin(\epsilon)}{\epsilon}\right)^2, \quad \left(\frac{\sin(\epsilon)}{\epsilon}\right)^2 \left(1 + \frac{1 - \cos(\epsilon)}{2}\right)$$

Hochbruch & Lubich (1999)

## SMC4 method

- To avoid resonance in simplified problem,

$$\phi(0) > 0,$$

$$\phi\left(\frac{\pi}{2} + 2m\pi\right) > 0,$$

$$\phi(\pi + 2m\pi) \leq 0, \quad m \in \mathbb{Z},$$

$$\text{If } \phi(\pi + 2m\pi) = 0 \text{ then } \phi'(\pi + 2m\pi) = 0.$$

- To conserve consistency of order 4 for the more general problem, with the assumptions of regularity,  $g$  Lipschitz in  $L^2$ -norm on its second variable and  $\|A^2u\|_2$  bounded,

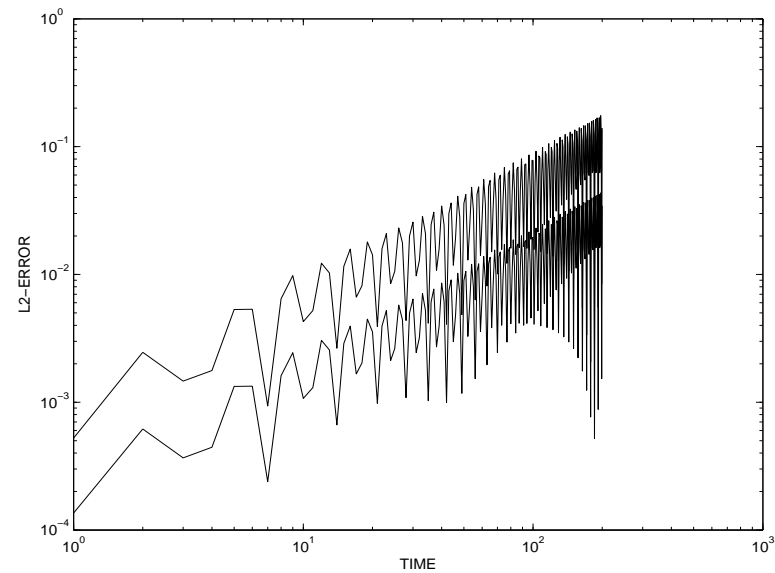
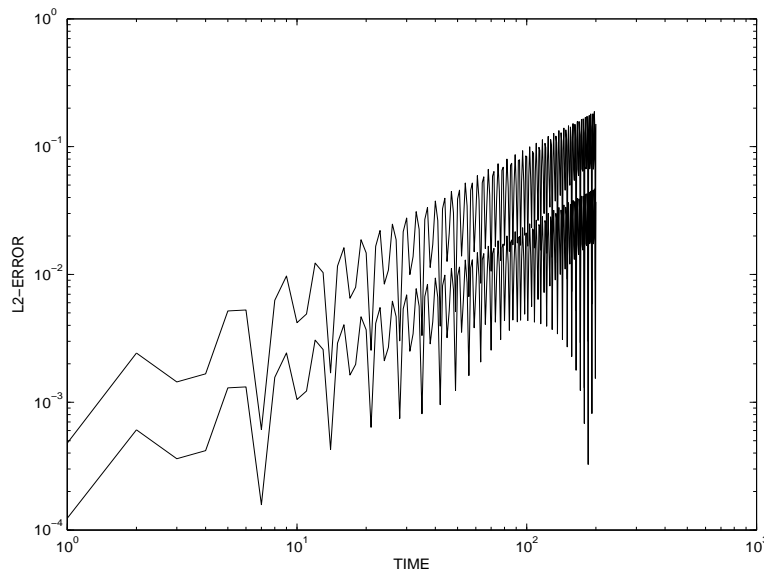
$$\frac{\phi(\epsilon) - 1}{\epsilon^4} \text{ bounded for real } \epsilon$$

$$\text{Possible filter function: } \left(\frac{\sin(\epsilon)}{\epsilon}\right)^2 \left[1 + \frac{2}{3}(1 - \cos(\epsilon))\right]$$

## Numerical comparison

Growth of error with time for the nonlinear wave equation

Gautschi ( $h = 1/10, 1/20$ ), pseudospectral ( $M = 500$ )



Annihilating frequencies

filter function

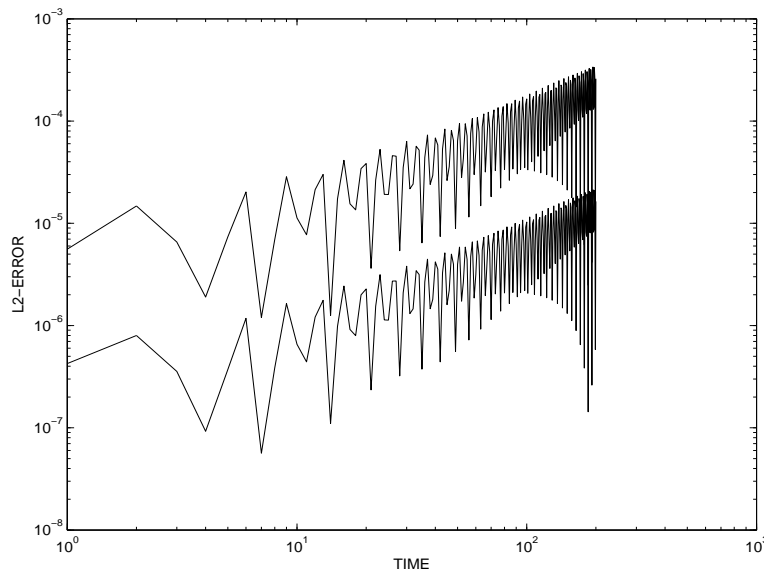
$$\phi(\epsilon) = \left(\frac{\sin(\epsilon)}{\epsilon}\right)^2 \left(1 + \frac{1 - \cos(\epsilon)}{2}\right)$$



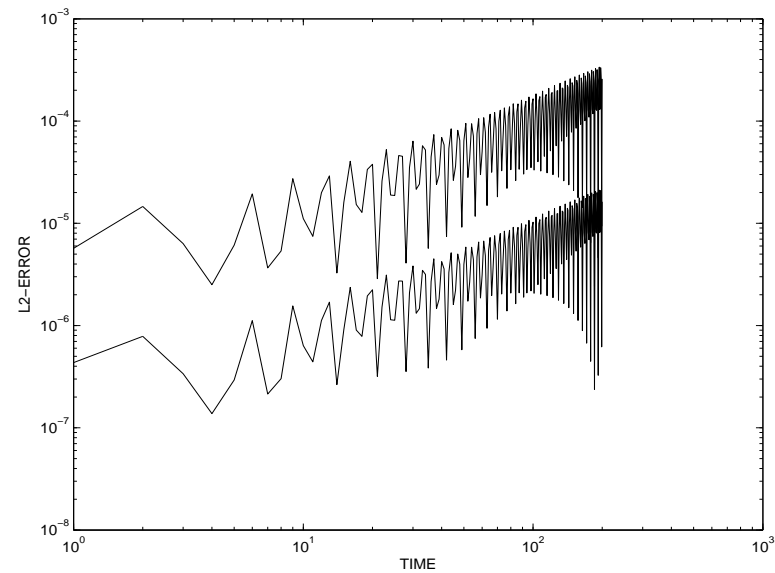
## Numerical comparison

Growth of error with time for the nonlinear wave equation

SMC4 ( $h = 1/10, 1/20$ ), pseudospectral ( $M = 500$ )



Annihilating frequencies



filter function

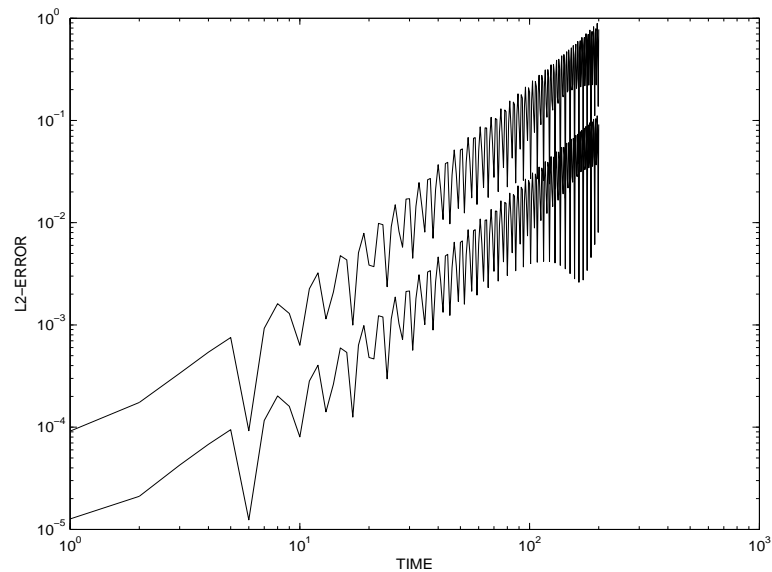
$$\phi(\epsilon) = \left(\frac{\sin(\epsilon)}{\epsilon}\right)^2 \left[1 + \frac{2}{3}(1 - \cos(\epsilon))\right]$$

## Qualitative properties of symmetric methods

Better behaviour with respect to growth of error with time

MC3 (not symmetric) ( $h = 1/10, 1/20$ ),

pseudospectral ( $M = 250$ )



Quadratic error growth while linear before

## Numerical integration of Euler-Bernoulli equation

$$\begin{aligned}u_{tt} &= -u_{xxxx} - 0,4u^3, & x \in [0, 2\pi], & 0 \leq t \leq T, \\u(x, 0) &= \frac{5}{2}e^{\frac{\cos(x)}{5}} - \frac{5}{2}, \\v(x, 0) &= e^{\frac{\sin(x)}{5}} - \frac{5}{2},\end{aligned}$$

periodic boundary conditions

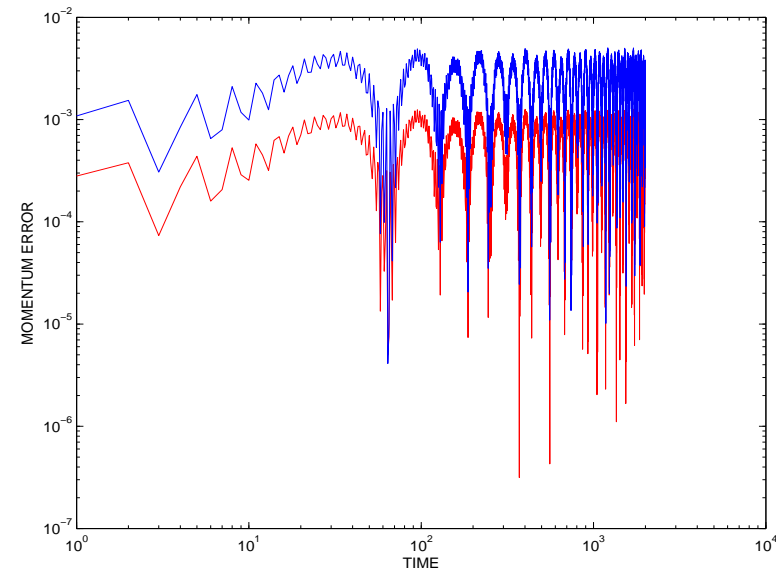
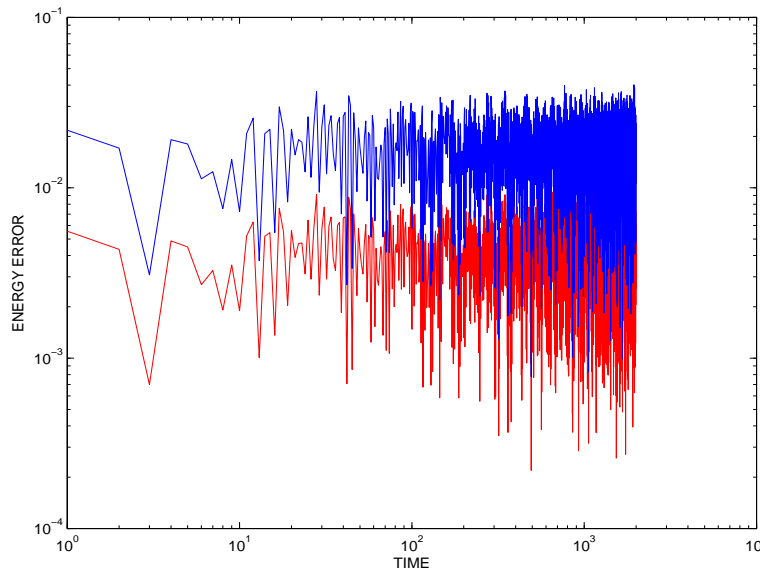
## Conserved quantities

$$\begin{aligned}H &= \int_a^b \left[ \frac{1}{2}u_t^2 + \frac{1}{2}u_{xx}^2 + 0,1u^4 \right] dx, \\M &= - \int_a^b u_t u_x dx.\end{aligned}$$

## Growth of energy and momentum error with time

Gautschi (filtered) ( $h = 1/20, 1/40$ ),

pseudospectral ( $M = 10$ )



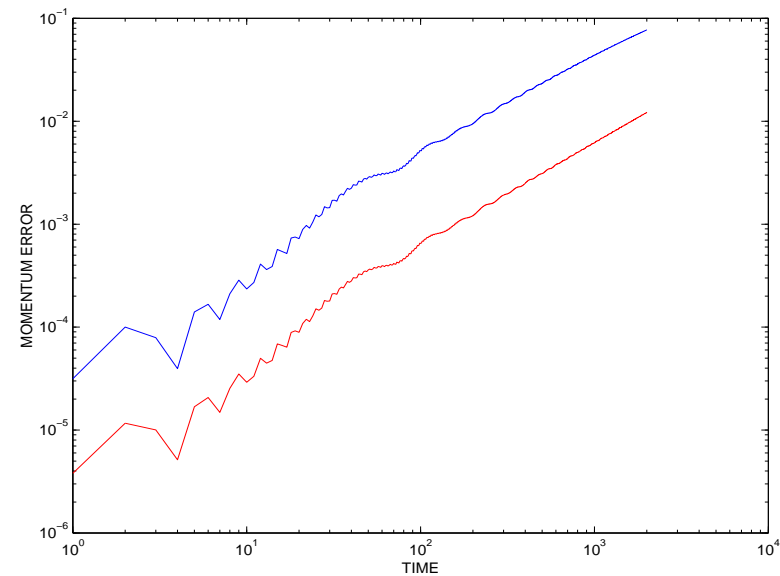
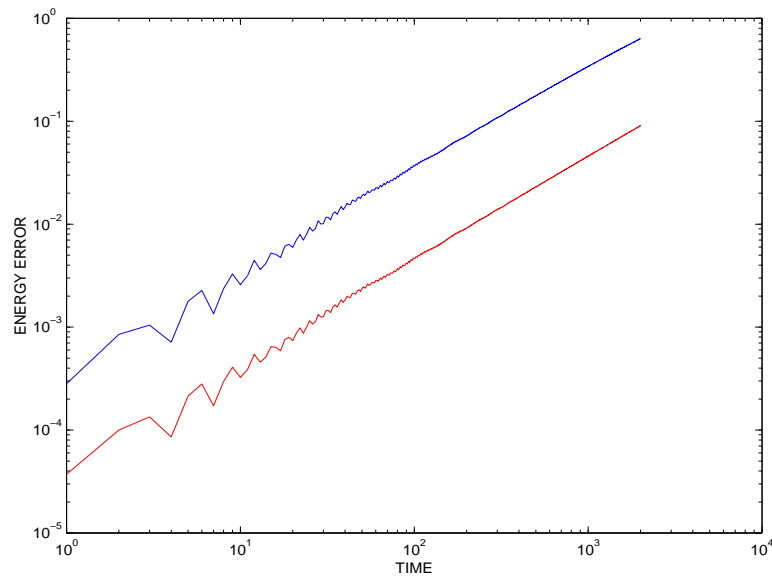
‘bounded ’

$$h = 1/20, \rightarrow \lambda_M h = 20, \quad h = 1/40, \rightarrow \lambda_M h = 10$$

## Growth of energy and momentum error with time

MC3 (filtered) ( $h = 1/20, 1/40$ ),

pseudospectral ( $M = 10$ )

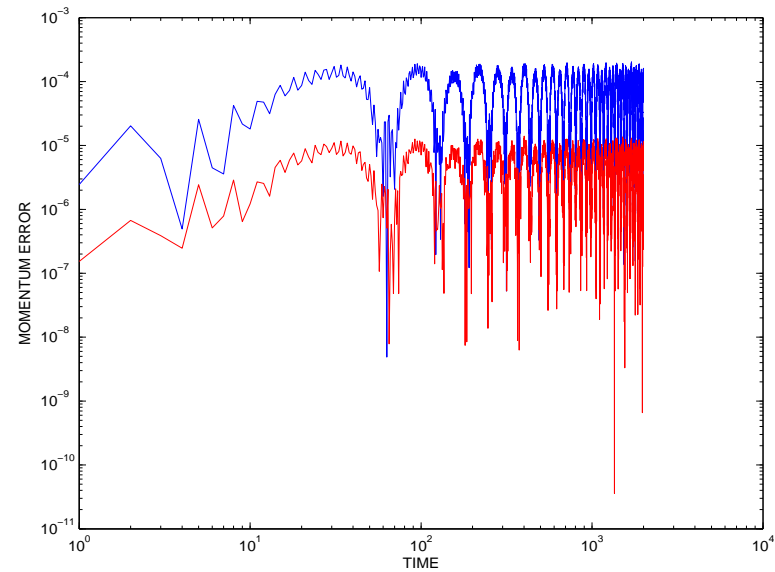
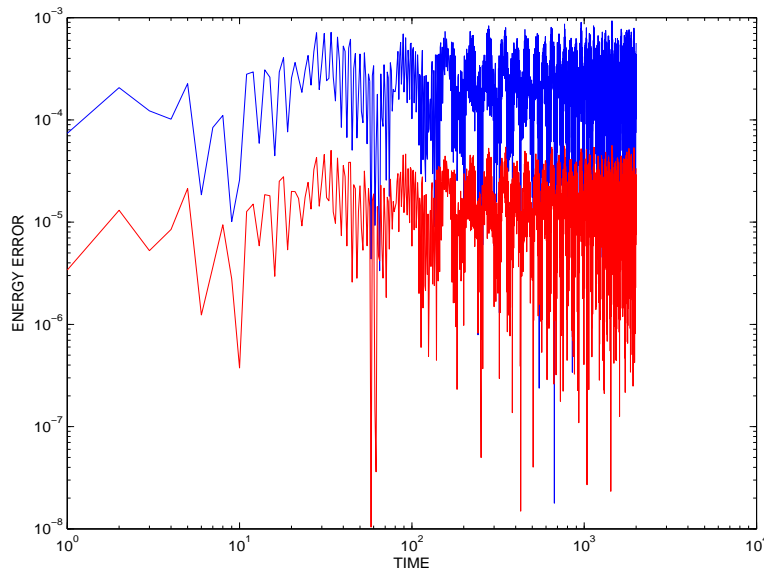


linear error growth

## Growth of energy and momentum error with time

SMC4 (filtered) ( $h = 1/20, 1/40$ ),

pseudospectral ( $M = 10$ )



'bounded'