

An overview of Celestial Mechanics: from the stability of the planets to flight dynamics



Alessandra Celletti



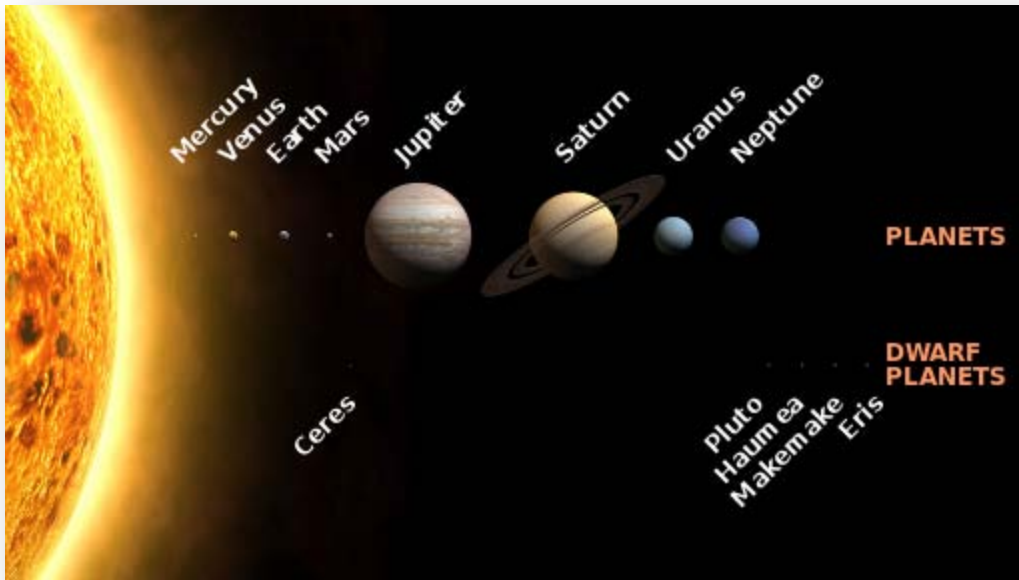
Dipartimento di Matematica
Università di Roma Tor Vergata
celletti@mat.uniroma2.it

Image: CICLOPS, JPL, ESA, NASA

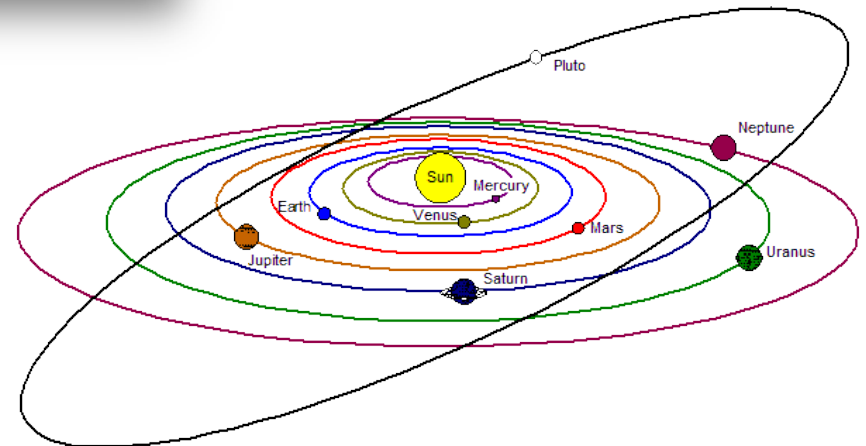
SUMMARY

- The Solar System
- Celestial Mechanics
- The 2-body problem
- The 3-body problem
- Chaos
- Orbital resonances
- Interplanetary highways
- Spin-orbit resonances
- Is the Solar system stable?

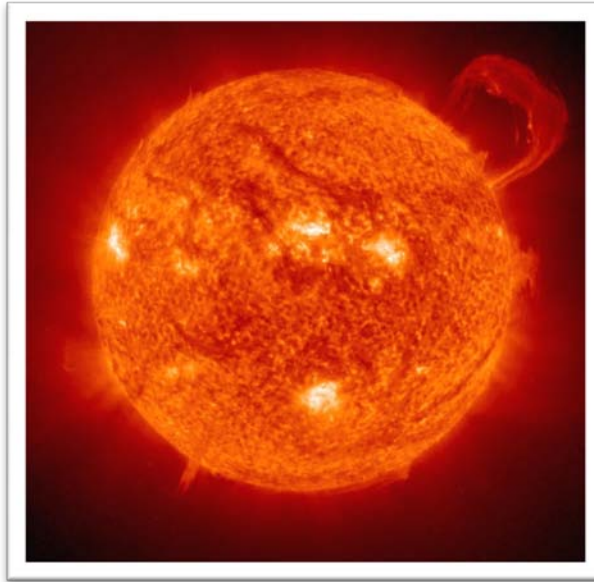
1. The solar system



- Sun
- Rocky planets
- Gas planets
- Dwarf planets
- Satellites
- Asteroids and comets

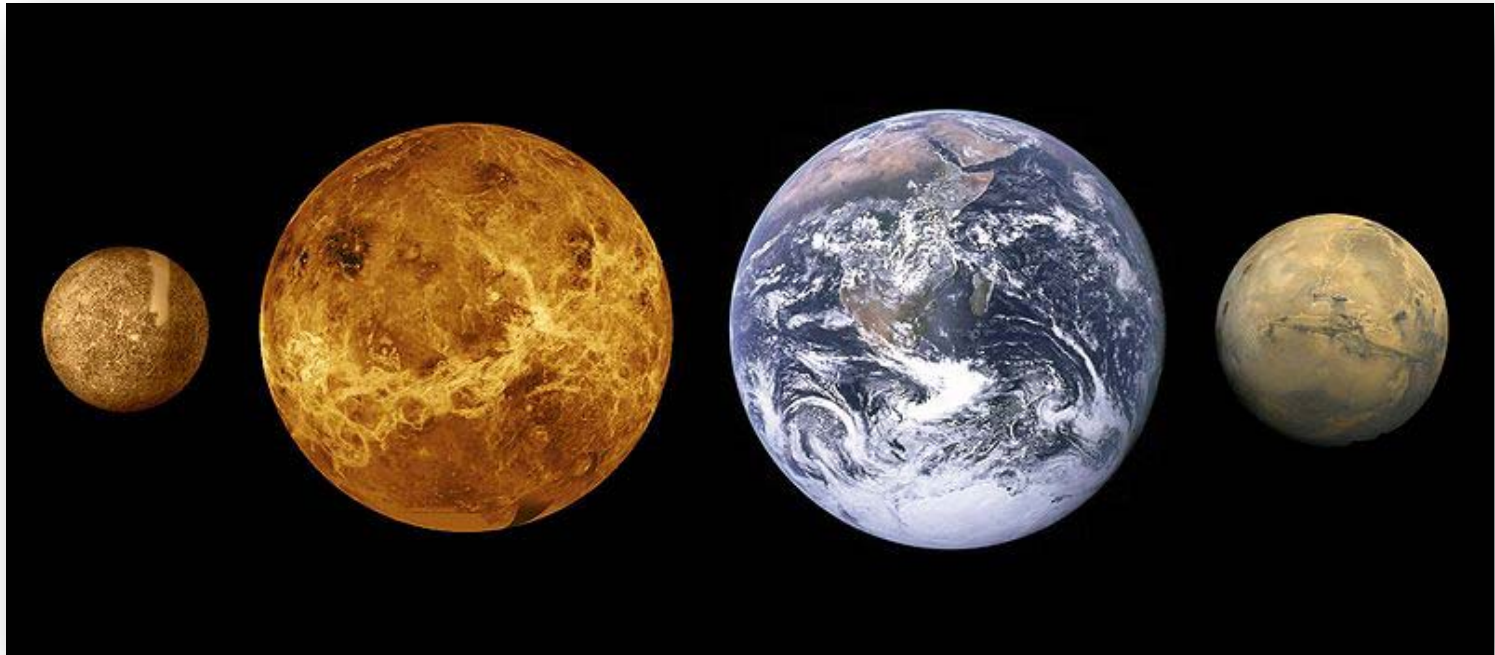


Il Sole



- Average size star at the border of a spiral arm of the Milky Way
- At about half of its life (4.5 billion years)
- Will evolve as a red giant and then a white dwarf
- Mass: $2 \cdot 10^{30}$ Kg
- Radius 695,000 Km
- Composition: H - He

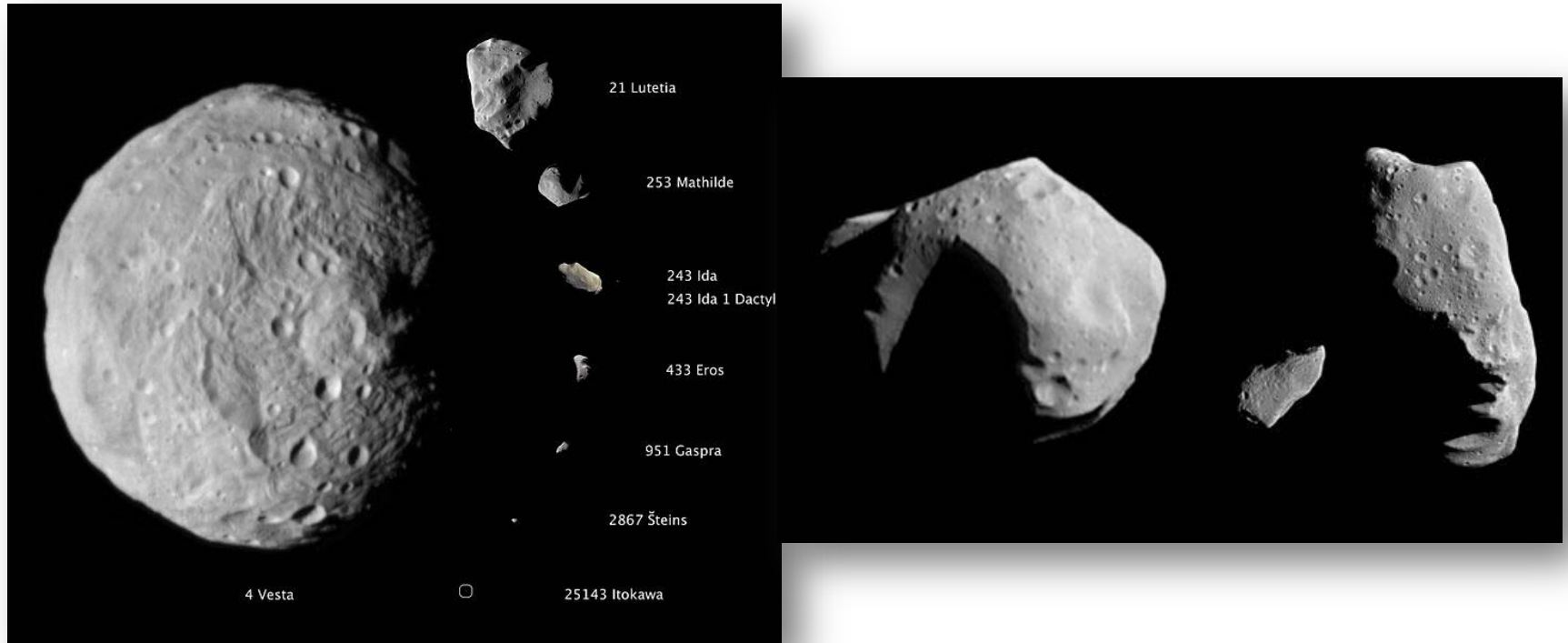
Internal rocky planets



Mercury, Venus, Earth, Mars

Small, rocky, none or few satellites

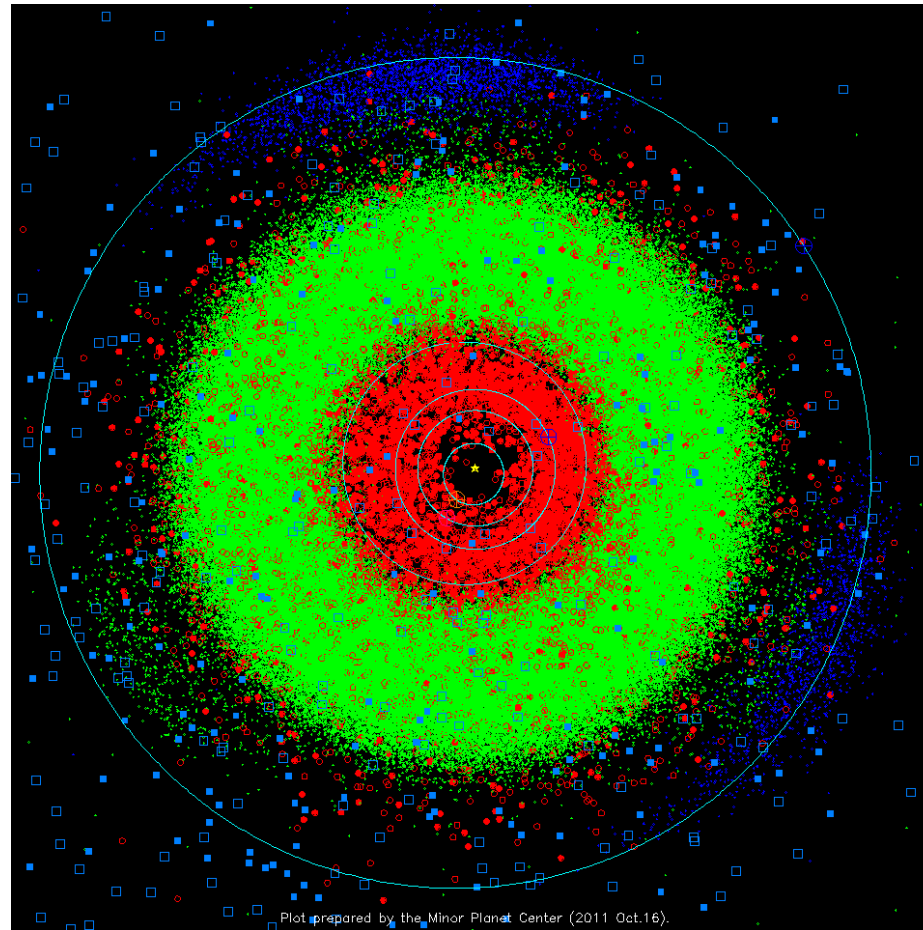
Asteroids: 580.000 oggetti catalogati



Between Mars and Jupiter, irregular forms and sizes, sometimes with satellites

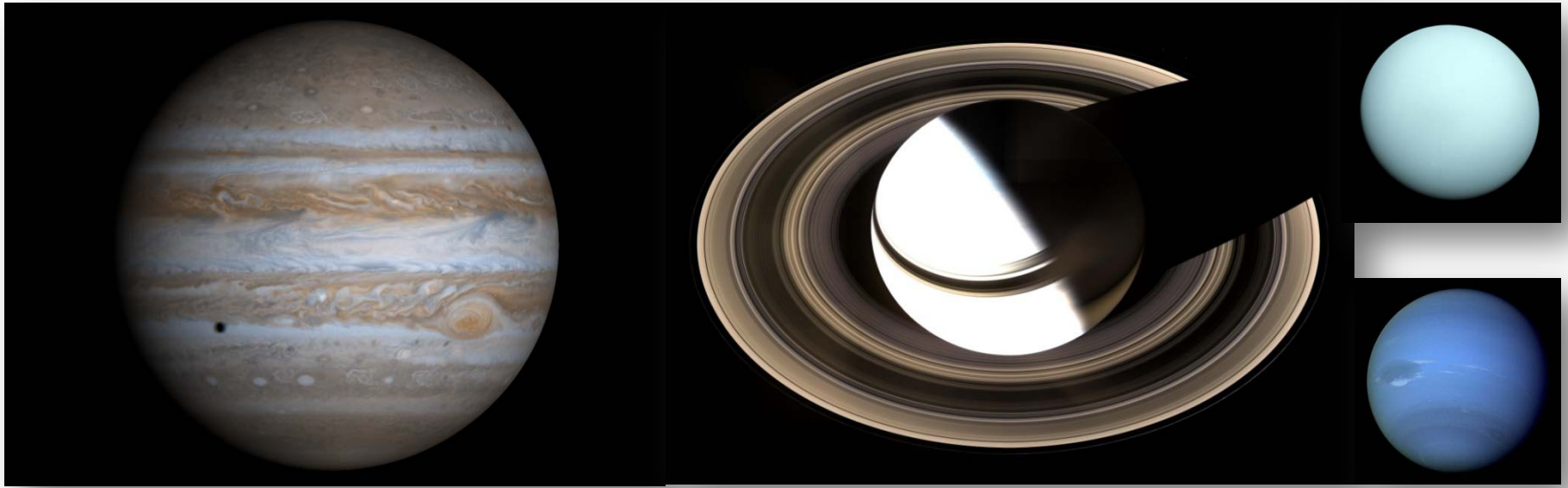


Internal Solar system



Asteroids-green, NEO-red, comets-blue.

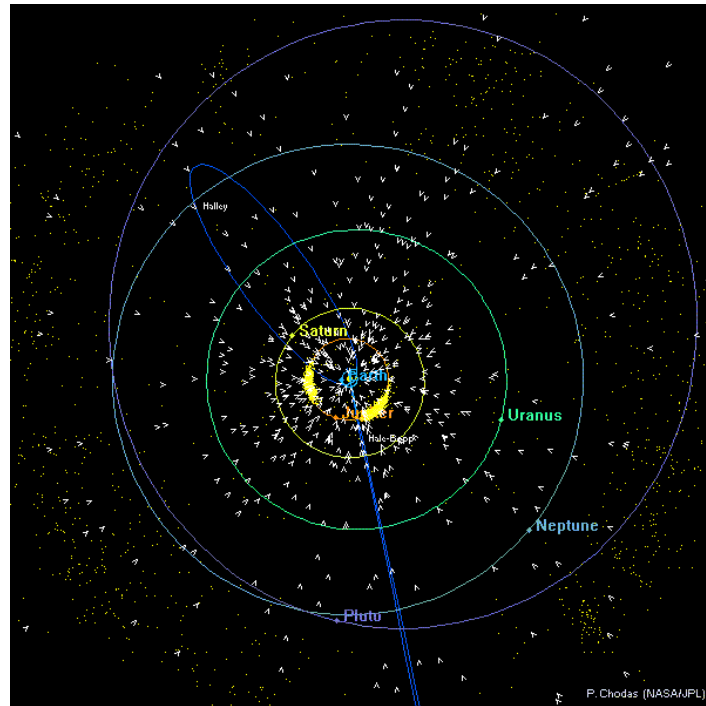
External gaseous planets



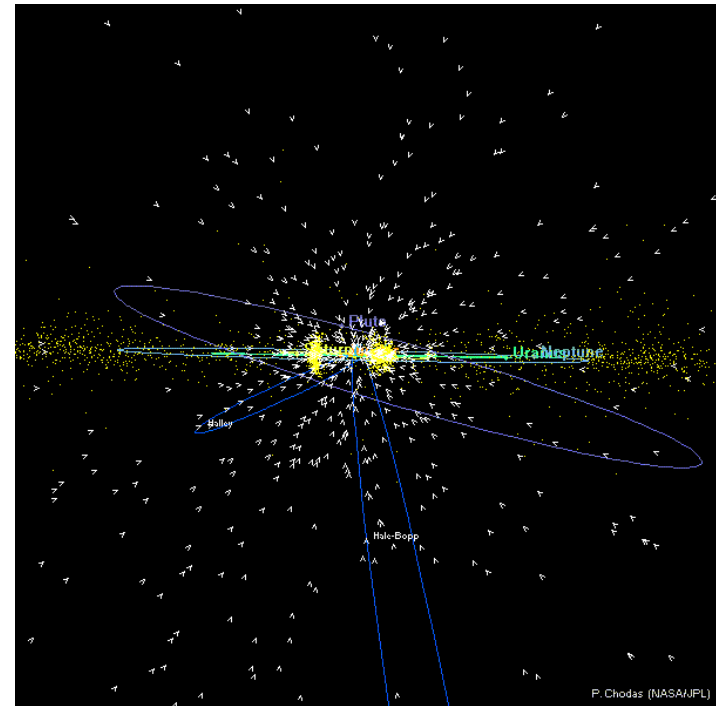
Jupiter, Saturn, Uranus, Neptune

Big, gas, many satellites and rings

External Solar system: asteroids-yellow, in comets-white

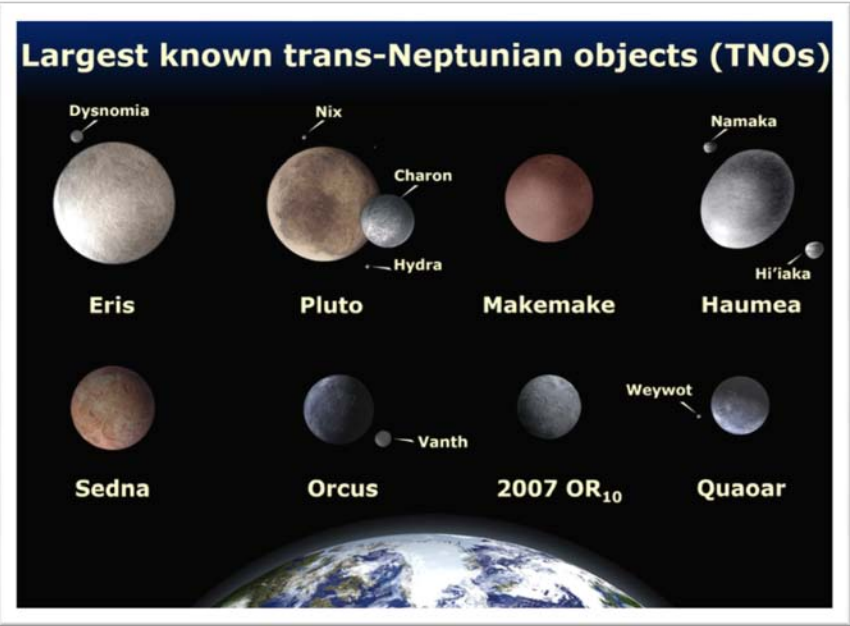
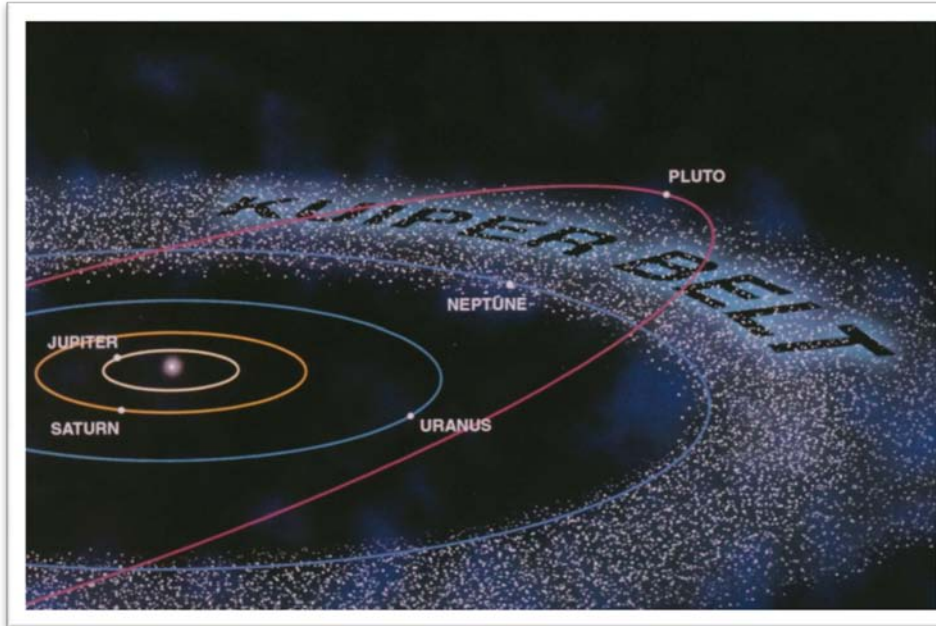


From above



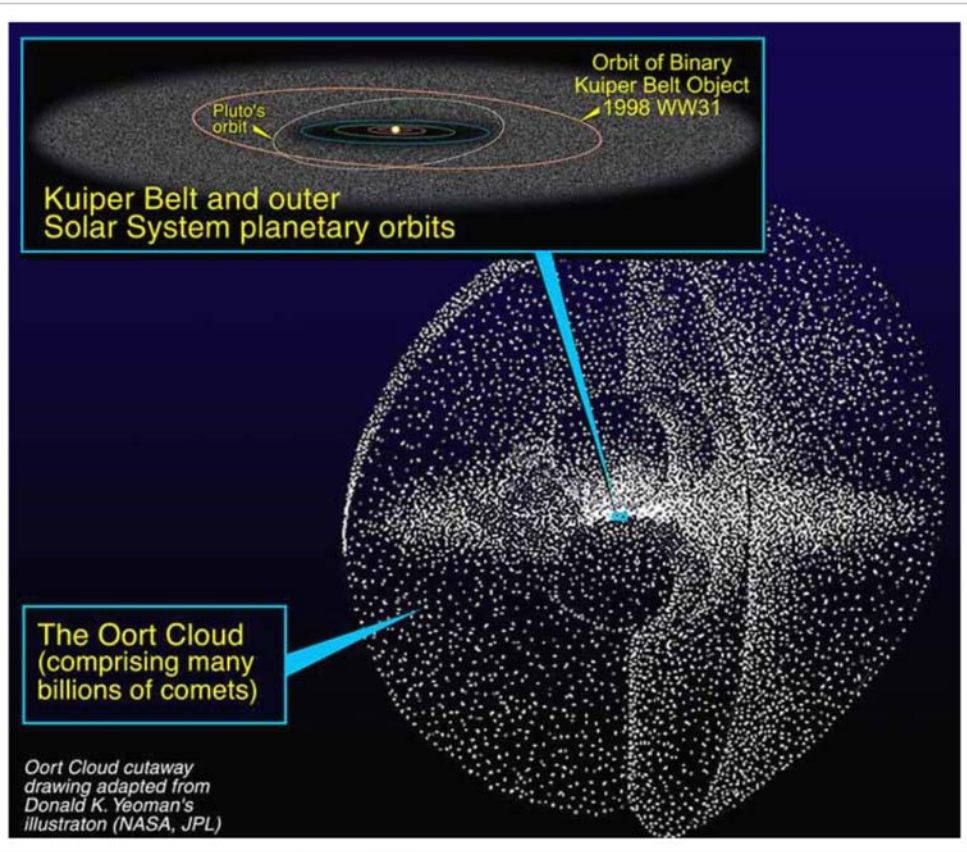
Profile

Kuiper belt



Thousands rocky-icy objects (among which Pluto), at the border of the Solar system

Oort cloud



- 30,000 - 100,000 AU
- Billion icy objects
- Long-period comets, inserted in the Solar system by strong perturbations (close encounter of a star or passage of the Sun through a giant molecular cloud).

1 AU = Sun-Earth distance = 150 million km.

2. Celestial Mechanics

- **CELESTIAL MECHANICS** studies the dynamics of the objects in the Solar system: planets, satellites, asteroids, etc.
- **CELESTIAL MECHANICS** studies also the dynamics of **extrasolar planetary systems**
- **FLIGHT DYNAMICS** studies the motion of artificial satellites and interplanetary highways (first **space mission**: Sputnik 1 on 4 October 1957)



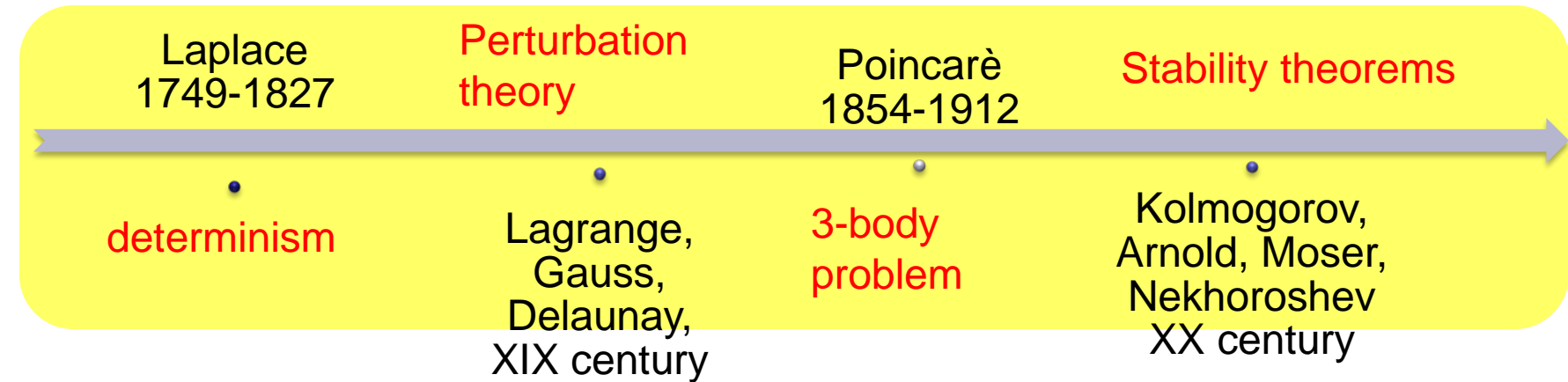
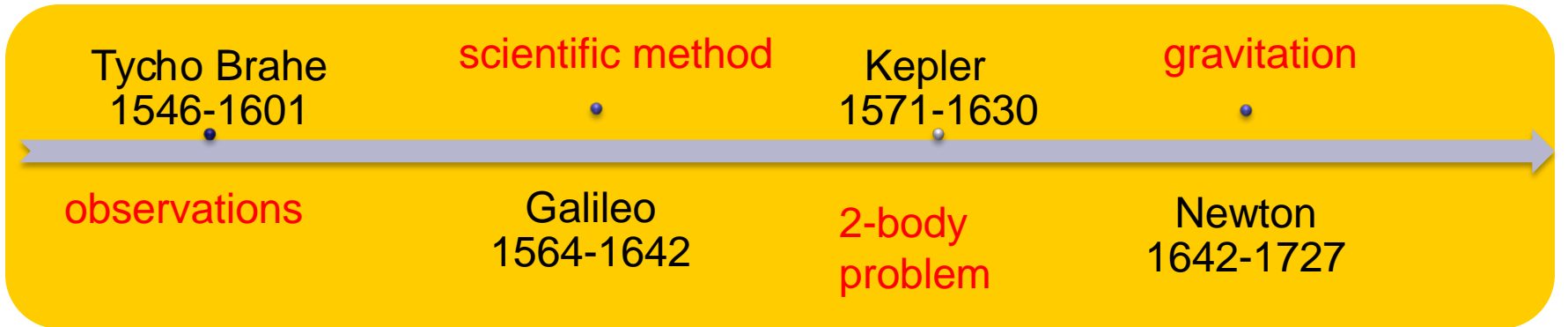
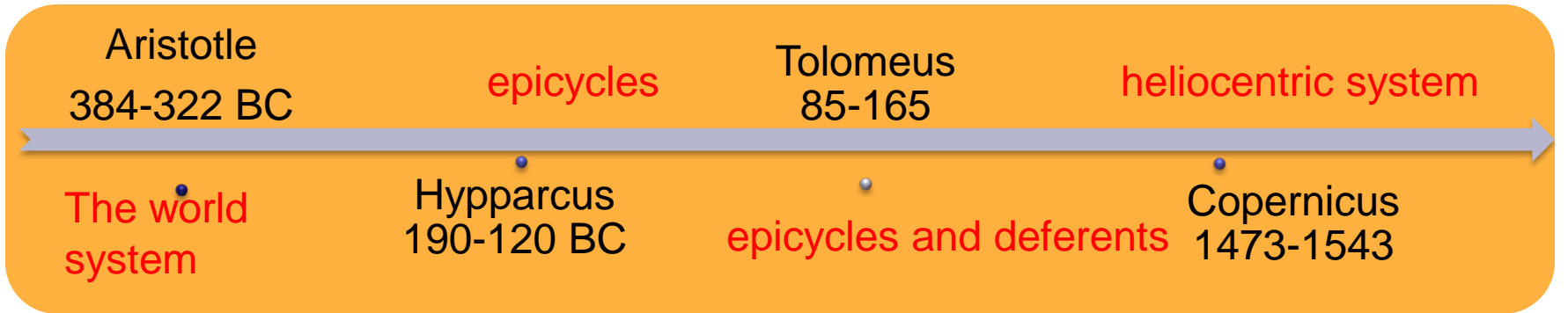
Annefrank

Wild 2

Tempel1



Immagini: NASA/JPL-Caltech/University of Maryland/Cornell, NASA and The Hubble Heritage Team (STScI/AURA), ESA-Hubble Collaboration, E. L. Wright (UCLA), The COBE Project, DIRBE, NASA , GALILEO/NASA/JPL



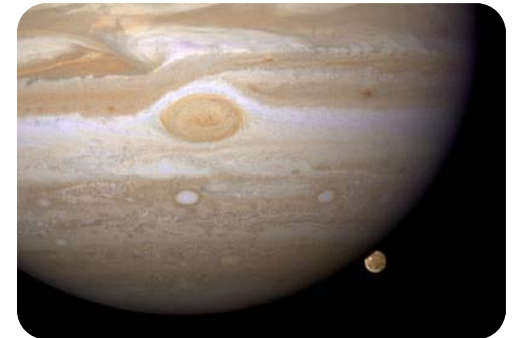
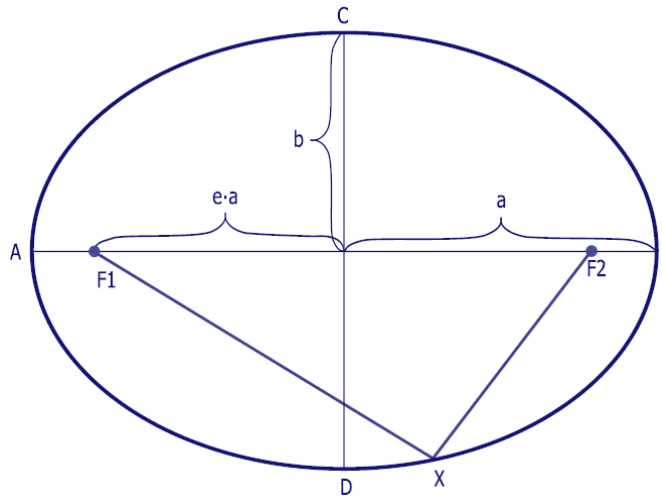
3. 2 body problem

- Simplified model considering only the interaction between **2** objects

- Newton's law:

$$F = - \frac{G M m}{d^2}$$

- Kepler's laws



Immagini: NASA, NASA/NEAR

Johannes Kepler (1571-1630)

- **Kepler** believed in the heliocentric theory of Copernicus
- He wrote several books where astronomy was mixed with mathematics, physics, philosophy and music
- Studied for several years the astronomical data on the motion of the planets, collected by **Tycho Brahe** (1546-1601), who built an astronomical observatory called "**Uraniborg**" – "The castle of the sky"
- Found 3 fundamental laws governing the 2-body problem

I Kepler law:

all planets move on ellipses with the Sun in one focus



kepler1.avi

II Kepler law: all planets sweep equal areas in equal times



kepler2a.avi



kepler2b.avi

III Kepler law:

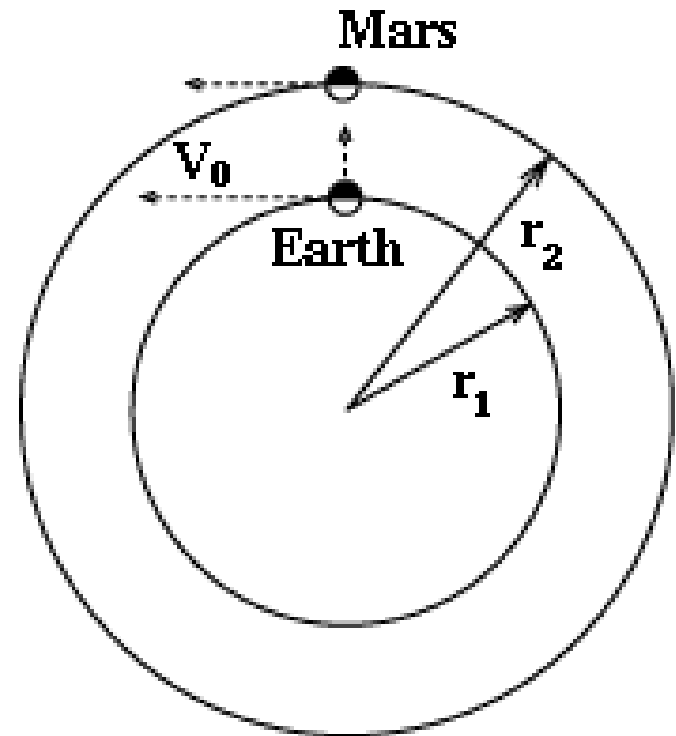
the square of the period of revolution is
proportional to the cube of the semimajor axis



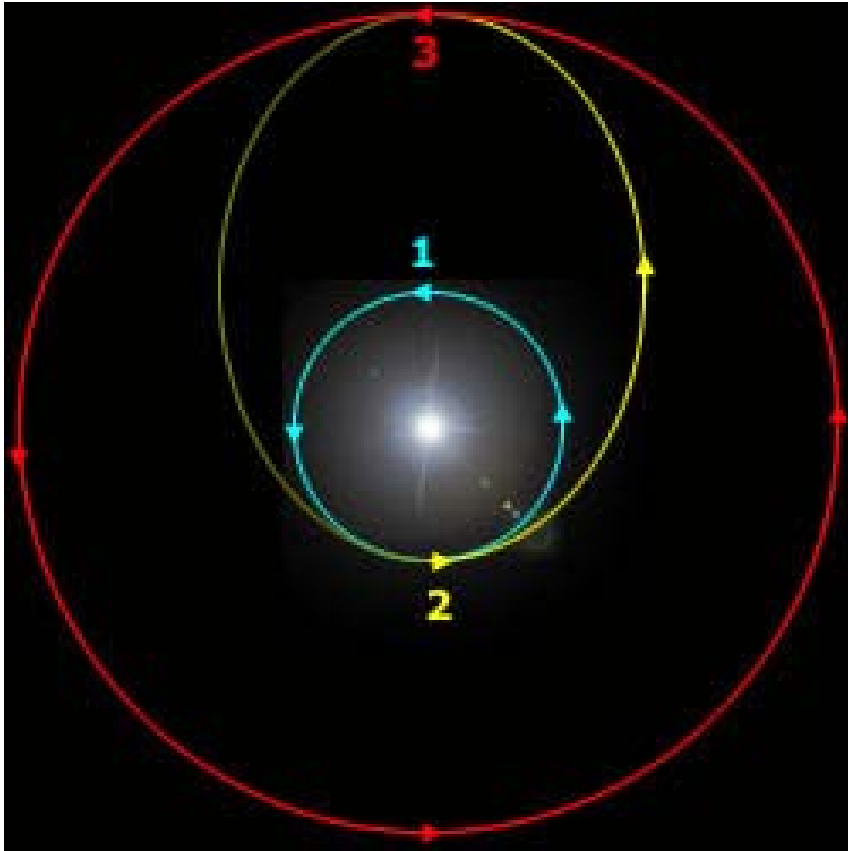
kepler3.avi

4. How NOT to go on Mars

- Earth and Mars on circular orbits with radii r_1 , r_2
- Wait for Earth-Mars conjunction and go on a straight line!
 - Gravity curves the trajectories
 - the orbit of Mars is reached perpendicularly
 - the Sun has a gravitazionali influence on the satellite.

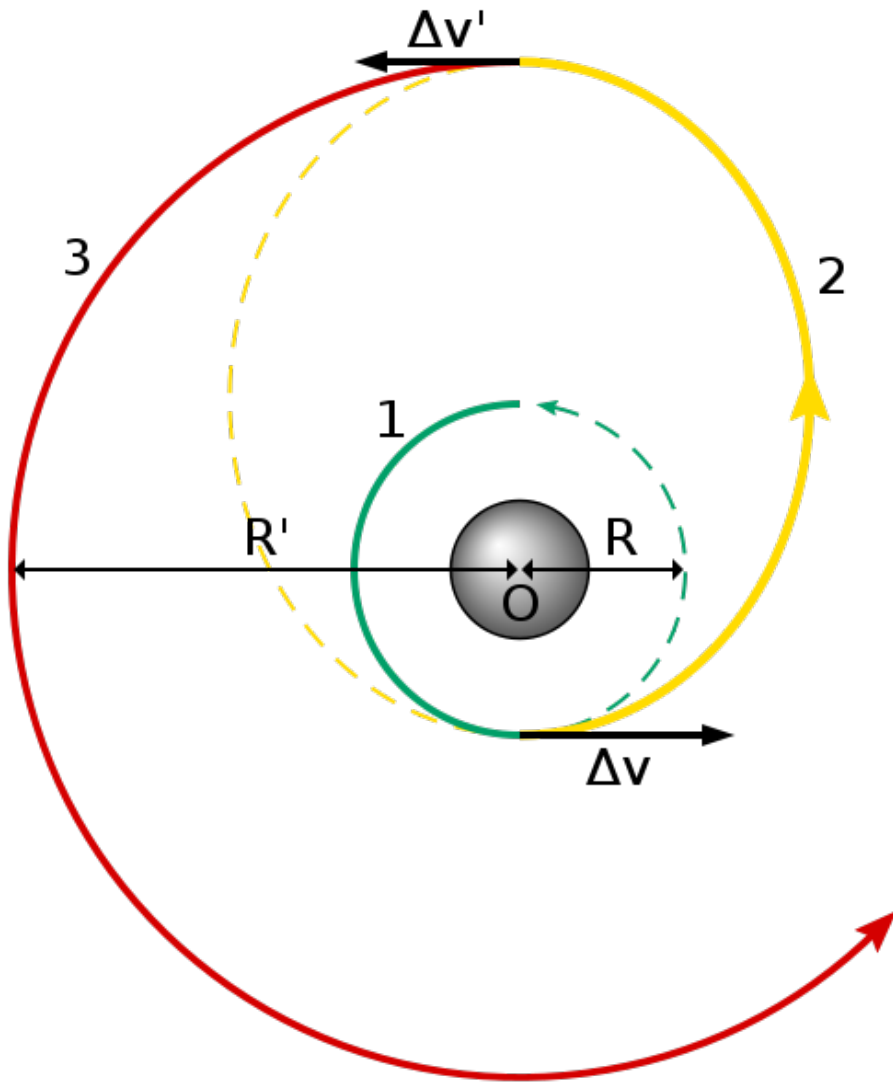


5. How to go on Mars



- Walter Hohmann (1880-1945) orbits
- 1 = initial orbit
- 2 = Hohmann transfer orbit
- 3 = target orbit

- Orbit 2 has perihelion on orbit 1 and aphelion on orbit 3
- Transfer with less fuel



- Switch the engines to insert the satellite in orbit 2 and then in orbit 3 (Δv)
- Δv measures the fuel consumption = cost of the mission
- *Launch window* is the time interval to have that the satellite reaches Mars

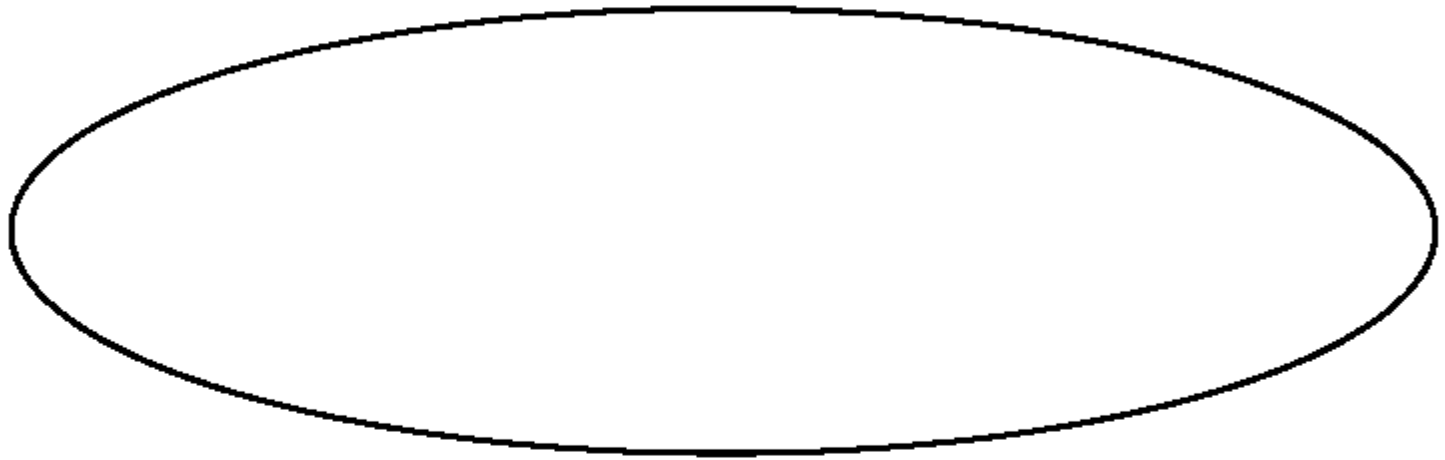
6. The three body problem

- What happens when we consider 3 bodies, e.g. **Sun-Earth-Jupiter** ?
- Kepler laws are only an approximation and the 3 body problem cannot be solved exactly!
- **Perturbation theory**: allows to compute *successive approximations* of the solution of the three body problem
- **Sun-Earth-Jupiter** : $\text{mass}(\textit{Jupiter}) = \text{mass}(\textit{Sun}) / 1000 \rightarrow$

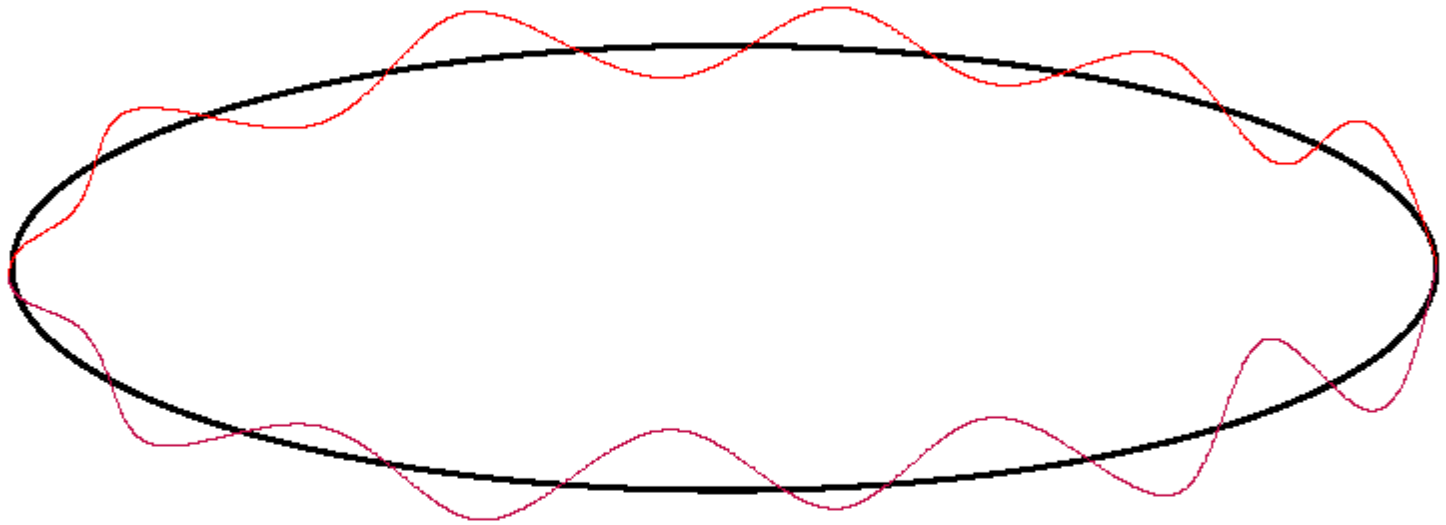
2-body Sun-Earth

+

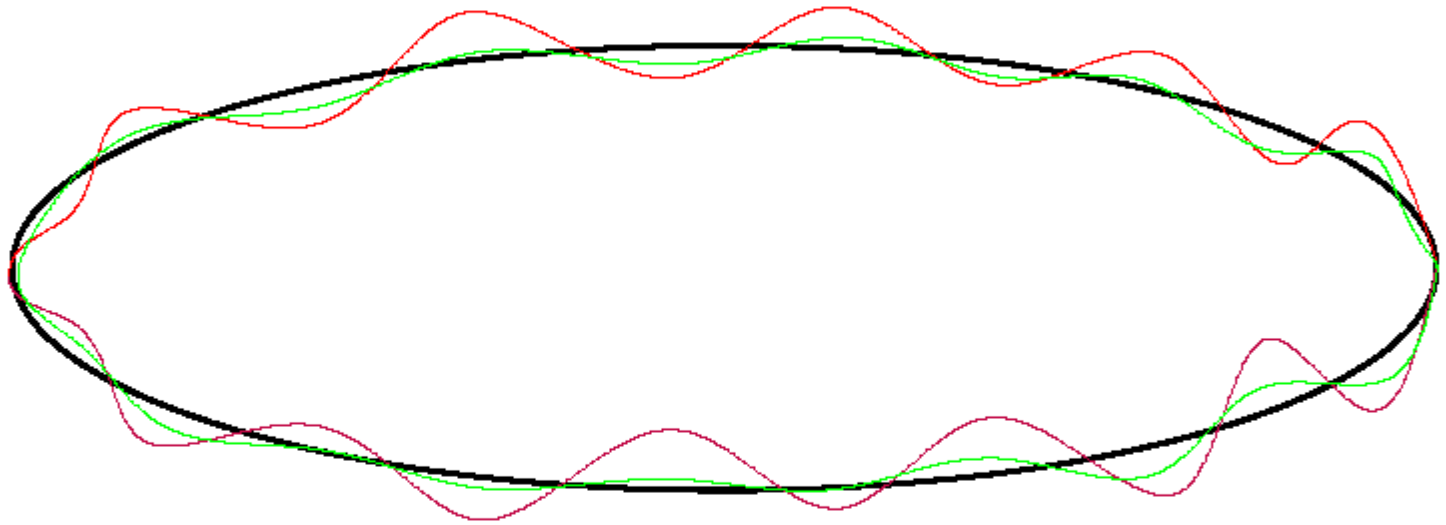
Small perturbation due to Jupiter



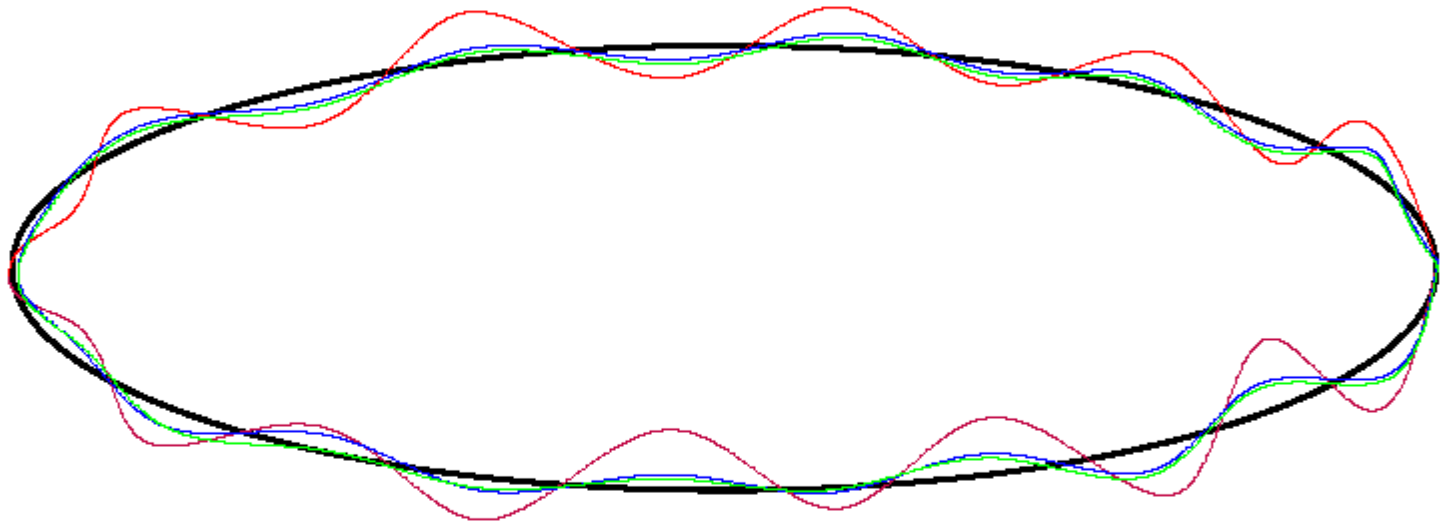
Keplerian ellipse: basic approximation



First approximation (red curve)



Second approximation (green curve)



Third approximation (blue curve)

- Perturbation theory allows to determined an approximate solution of the equations of motion (Laplace, Lagrange, Delaunay, Leverrier, etc., XVIII-XIX century).
- Charles Delaunay (1816-1872) developed a very precise lunar motion based on perturbation theory.



THÉORIE

MOUVEMENT DE LA LUNE.

CHAPITRE PREMIER.

EQUATIONS DIFFÉRENTIELLES DU MOUVEMENT DE LA LUNE. — MOUVEMENT ELLIPTIQUE. — VARIATION DES CONSTANTES DU MOUVEMENT ELLIPTIQUE.

1. Soient X, Y, Z les coordonnées de la Terre rapportées à des axes rectangulaires fixes dans l'espace; ξ, η, ζ les coordonnées de la Lune, et ξ', η', ζ' celles du Soleil rapportées aux mêmes axes; M la masse de la Terre, m celle de la Lune, et m' celle du Soleil.

Le Soleil, la Lune et la Terre étant supposés s'attirer mutuellement d'après la loi de Newton, les équations différentielles du mouvement de la Terre seront

$$\frac{d^2 X}{dt^2} = \frac{m'(\xi - X)}{[(1-X)^2 + (1-Y)^2 + (1-Z)^2]^{3/2}} - \frac{m^2(\xi - X)}{[(\xi - X)^2 + (\eta - Y)^2 + (\zeta - Z)^2]^{3/2}}$$

$$\frac{d^2 Y}{dt^2} = \frac{m'(\eta - Y)}{[(1-X)^2 + (1-Y)^2 + (1-Z)^2]^{3/2}} - \frac{m^2(\eta - Y)}{[(\xi - X)^2 + (\eta - Y)^2 + (\zeta - Z)^2]^{3/2}}$$

$$\frac{d^2 Z}{dt^2} = \frac{m'(\zeta - Z)}{[(1-X)^2 + (1-Y)^2 + (1-Z)^2]^{3/2}} - \frac{m^2(\zeta - Z)}{[(\xi - X)^2 + (\eta - Y)^2 + (\zeta - Z)^2]^{3/2}}$$

T. XXVIII.

$$\begin{aligned} \frac{d^2}{dt^2} \cos(x - \nu') &= \left(1 + 2e^2 + \frac{23e^4}{16}\right) \cos(x - g' - f') \\ &+ \left(3e' + \frac{11}{4}e'^2\right) \cos(x - g' - 3f') \\ &+ \left(e' + \frac{5}{2}e'^2\right) \cos(x - g') \\ &+ \left(\frac{53}{8}e'^2 + \frac{31}{16}e'^4\right) \cos(x - g' - 3f') \\ &+ \left(\frac{11}{8}e'^2 + \frac{49}{16}e'^4\right) \cos(x - g' + f') \\ &+ \frac{27}{8}e'^2 \cos(x - g' - 4f') \\ &+ \frac{23}{12}e'^2 \cos(x - g' + 2f') \\ &+ \frac{2655}{128}e'^2 \cos(x - g' - 5f') \\ &+ \frac{343}{128}e'^2 \cos(x - g' + 3f'); \end{aligned}$$

$$\begin{aligned} \frac{d^2}{dt^2} \cos(x - 3\nu') &= \left(1 - 6e^2 + \frac{43}{16}e^4\right) \cos(x - 3g' - 3f') \\ &+ (5e' - 22e'^2) \cos(x - 3g' - 4f') \\ &- \left(e' - \frac{5}{4}e'^2\right) \cos(x - 3g' - 2f') \\ &+ \left(\frac{127}{8}e'^2 - \frac{3665}{48}e'^4\right) \cos(x - 3g' - 5f') \\ &+ \left(\frac{1}{8}e'^2 + \frac{1}{48}e'^4\right) \cos(x - 3g' - f') \\ &+ \frac{163}{4}e'^2 \cos(x - 3g' - 6f') \\ &+ \frac{35413}{384}e'^2 \cos(x - 3g' - 7f') \\ &+ \frac{1}{384}e'^2 \cos(x - 3g' + f'); \end{aligned}$$

(*) La valeur de $\frac{d^2}{dt^2} \cos(x - 3\nu')$, calculée jusqu'aux quantités du quatrième ordre par rapport à e' , ne renferme aucun terme en $\cos x - 3g'$.

“Theorie du Mouvement de la Lune”
C. Delaunay

Preliminary computations

auquel on aurait dû s'arrêter, d'après ce qui vient d'être dit, et cela pour des raisons spéciales qui seront indiquées plus tard (chapitre IV).

Ajoutons encore que, e' étant environ trois fois plus petit que γ et e , dans le rejet des termes d'un ordre supérieur à celui auquel on voulait s'arrêter, on a regardé e'' , e''' , $e^{(4)}$, comme des quantités des quatrième, cinquième, sixième ordres; $e^{(4)}$ comme une quantité du huitième ordre, etc.

En opérant conformément aux explications qui précèdent, on a trouvé pour R la valeur suivante :

$$\begin{aligned}
 R = & \frac{e}{2\alpha} \\
 & + \frac{\pi e^2}{2\alpha^2} \left[\frac{1}{4} \gamma^2 e - \frac{1}{8} \gamma^2 e^2 - \frac{1}{8} \gamma^2 e^2 + \frac{1}{4} \gamma^2 e^2 - \frac{1}{2} \gamma^2 e^2 - \frac{1}{2} \gamma^2 e^2 + \frac{15}{16} e^2 e^2 \right. \\
 & + \frac{1}{2} \gamma^2 e^2 + \frac{1}{2} \gamma^2 e^2 - \frac{1}{2} \gamma^2 e^2 e^2 - \frac{15}{16} \gamma^2 e^2 e^2 - \frac{15}{64} \gamma^2 e^2 e^2 \\
 & + \left(\frac{2}{63} - \frac{45}{16} \gamma^2 + \frac{45}{63} \gamma^2 + \frac{45}{63} \gamma^2 \right) \frac{e^2}{\alpha^2} \\
 & + \left[\frac{1}{4} - \frac{3}{2} \gamma^2 - \frac{15}{8} \gamma^2 - \frac{15}{8} \gamma^2 e^2 + \frac{1}{2} \gamma^2 + \frac{15}{2} \gamma^2 e^2 + \frac{15}{2} \gamma^2 e^2 + \frac{69}{64} \gamma^2 e^2 + \frac{75}{64} \gamma^2 e^2 \right. \\
 & + \frac{39}{64} \gamma^2 e^2 - \frac{15}{8} \gamma^2 e^2 - \frac{15}{8} \gamma^2 e^2 - \frac{69}{32} \gamma^2 e^2 - \frac{75}{8} \gamma^2 e^2 e^2 - \frac{65}{32} \gamma^2 e^2 - \frac{145}{128} \gamma^2 e^2 \\
 & \left. - \left(\frac{5}{16} - 5\gamma^2 + \frac{5}{16} \gamma^2 + \frac{5}{16} \gamma^2 \right) \frac{e^2}{\alpha^2} \right] \cos(\alpha h + \alpha g + \alpha l - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[-\frac{1}{2} \gamma^2 e + 3\gamma^2 e + \frac{1}{16} \gamma^2 e - \frac{1}{4} \gamma^2 e - 3\gamma^2 e - \frac{1}{4} \gamma^2 e^2 + \frac{1}{2} \gamma^2 e^2 - \frac{1}{32} \gamma^2 e^2 \right. \\
 & \left. + \frac{3}{32} \gamma^2 e^2 - \frac{15}{16} \gamma^2 e^2 - \frac{9}{16} \gamma^2 e^2 \right] \cos t \\
 & + \left[\frac{1}{4} \gamma^2 e - \frac{1}{2} \gamma^2 e^2 - \frac{1}{2} \gamma^2 e^2 + \frac{15}{16} \gamma^2 e^2 - \frac{1}{2} \gamma^2 e^2 - \frac{15}{16} \gamma^2 e^2 e^2 - \frac{15}{16} \gamma^2 e^2 e^2 - \frac{15}{64} \gamma^2 e^2 e^2 \right. \\
 & \left. + \frac{161}{256} \gamma^2 e^2 + \frac{37}{4} \gamma^2 e^2 e^2 + \left(\frac{45}{63} \gamma^2 - \frac{315}{16} \gamma^2 e^2 + \frac{315}{63} \gamma^2 e^2 \right) \frac{e^2}{\alpha^2} \right] \cos t
 \end{aligned}$$

T. XXVIII.

$$\begin{aligned}
 & + \left[\frac{31}{8} \gamma^2 e^2 - \frac{31}{4} \gamma^2 e^2 e^2 - \frac{69}{64} \gamma^2 e^2 e^2 - \frac{115}{8} \gamma^2 e^2 e^2 \right] \cos(\alpha h + \alpha g + \alpha l - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[-\frac{153}{8} \gamma^2 e^2 + \frac{153}{4} \gamma^2 e^2 e^2 + \frac{663}{64} \gamma^2 e^2 e^2 + \frac{245}{8} \gamma^2 e^2 e^2 \right] \cos(\alpha h + \alpha g + \alpha l - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[\frac{815}{64} \gamma^2 e^2 - \frac{815}{32} \gamma^2 e^2 e^2 - \frac{4355}{128} \gamma^2 e^2 e^2 - \frac{3555}{1024} \gamma^2 e^2 e^2 \right] \cos(\alpha h + \alpha g + \alpha l - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[\frac{1}{63} \gamma^2 e^2 - \frac{1}{32} \gamma^2 e^2 e^2 - \frac{5}{128} \gamma^2 e^2 e^2 + \frac{11}{1024} \gamma^2 e^2 e^2 \right] \cos(\alpha h + \alpha g + \alpha l - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[\frac{3}{2} \gamma^2 e - \frac{1}{2} \gamma^2 e - \frac{5}{16} \gamma^2 e^2 + \frac{1}{4} \gamma^2 e^2 \right] \cos(\alpha g + \alpha l) \\
 & + \left[-\frac{2}{4} \gamma^2 e + \frac{2}{4} \gamma^2 e + \frac{39}{16} \gamma^2 e^2 - \frac{37}{4} \gamma^2 e^2 e^2 - \frac{29}{16} \gamma^2 e^2 \right. \\
 & \left. + \frac{5}{128} \gamma^2 e^2 + \frac{112}{32} \gamma^2 e^2 e^2 \right] \cos(\alpha g + \alpha l) \\
 & + \left[\frac{3}{2} \gamma^2 e - \frac{1}{2} \gamma^2 e - \frac{45}{8} \gamma^2 e^2 e^2 + \frac{81}{32} \gamma^2 e^2 e^2 \right] \cos(\alpha g + \alpha l + \alpha l') \\
 & + \left[\frac{3}{2} \gamma^2 e - \frac{1}{2} \gamma^2 e - \frac{45}{8} \gamma^2 e^2 e^2 + \frac{81}{32} \gamma^2 e^2 e^2 \right] \cos(\alpha g + \alpha l - \alpha l') \\
 & + \left[-\frac{1}{2} \gamma^2 e + \frac{3}{4} \gamma^2 e + \frac{3}{16} \gamma^2 e^2 + \frac{15}{4} \gamma^2 e^2 e^2 - \frac{3}{16} \gamma^2 e^2 \right. \\
 & \left. - \frac{1}{128} \gamma^2 e^2 - \frac{15}{32} \gamma^2 e^2 e^2 \right] \cos(\alpha h + \alpha l - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[-\frac{1}{2} \gamma^2 e + \frac{3}{4} \gamma^2 e + \frac{3}{16} \gamma^2 e^2 + \frac{15}{4} \gamma^2 e^2 e^2 \right] \cos(\alpha h - \alpha l - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[\frac{31}{4} \gamma^2 e - \frac{31}{4} \gamma^2 e + \frac{63}{8} \gamma^2 e^2 e^2 - \frac{369}{32} \gamma^2 e^2 e^2 - \frac{63}{8} \gamma^2 e^2 e^2 \right] \cos(\alpha h - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[-\frac{3}{4} \gamma^2 e + \frac{3}{4} \gamma^2 e - \frac{1}{8} \gamma^2 e^2 e^2 + \frac{3}{16} \gamma^2 e^2 e^2 - \frac{1}{8} \gamma^2 e^2 e^2 \right] \cos(\alpha h - \alpha h' - \alpha g' - \alpha l') \\
 & + \left[-\frac{1}{32} \gamma^2 e + \frac{1}{4} \gamma^2 e + \frac{1}{16} \gamma^2 e^2 - \frac{1}{16} \gamma^2 e^2 \right] \cos \alpha l
 \end{aligned}$$

(4, 0)

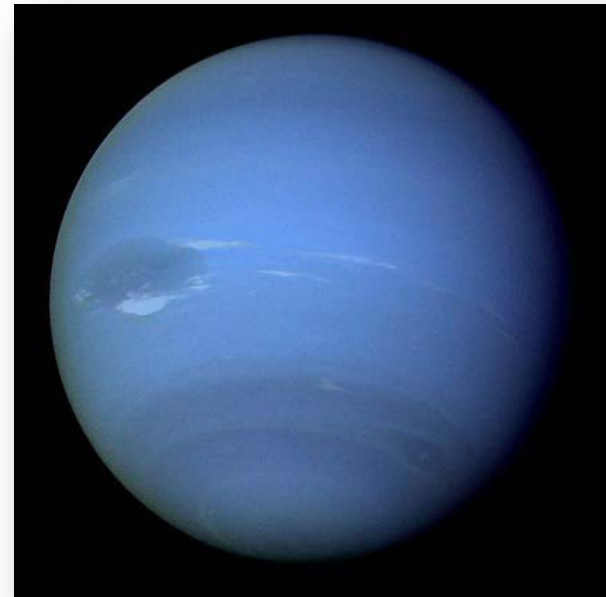
$$\begin{aligned}
 & - \frac{189}{64} e^2 \cos^2 \lambda h + \lambda g + 6l - \lambda h - \lambda g - 3l \\
 & - \frac{37}{64} e^2 \cos^2 \lambda h + \lambda g + 6l - \lambda h - \lambda g - 3l \\
 & - \frac{31}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{3}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & + \frac{193}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{119}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & + \frac{813}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & + \frac{1}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{1245}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{5}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{159}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{1}{32} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{1791}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{1}{32} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & + \frac{19417}{320} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{111}{5120} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & + \frac{15}{16} e^2 \cos^2 \lambda h + \lambda g + 3l
 \end{aligned}$$

T XXVIII.

$$\begin{aligned}
 & - \frac{35}{12} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{105}{32} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{135}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{105}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & + \frac{105}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{117}{32} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{135}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{105}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{1305}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{111}{64} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & + \frac{105}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{33}{128} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{35}{16} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & - \frac{35}{16} e^2 \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l \\
 & = m \frac{e^2}{128} \cos^2 \lambda h + \lambda g + 3l - \lambda h - \lambda g - 3l.
 \end{aligned}$$

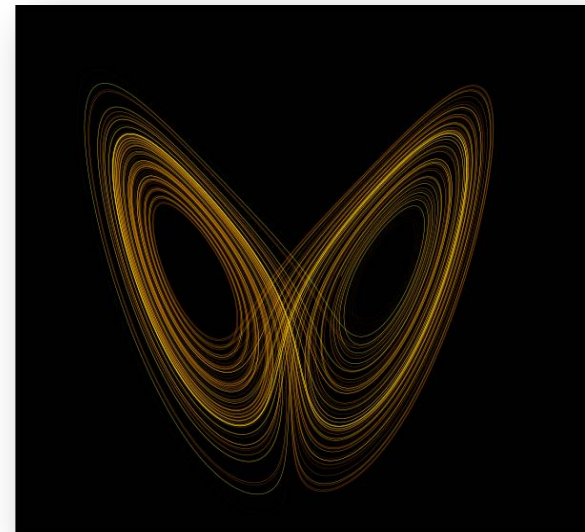
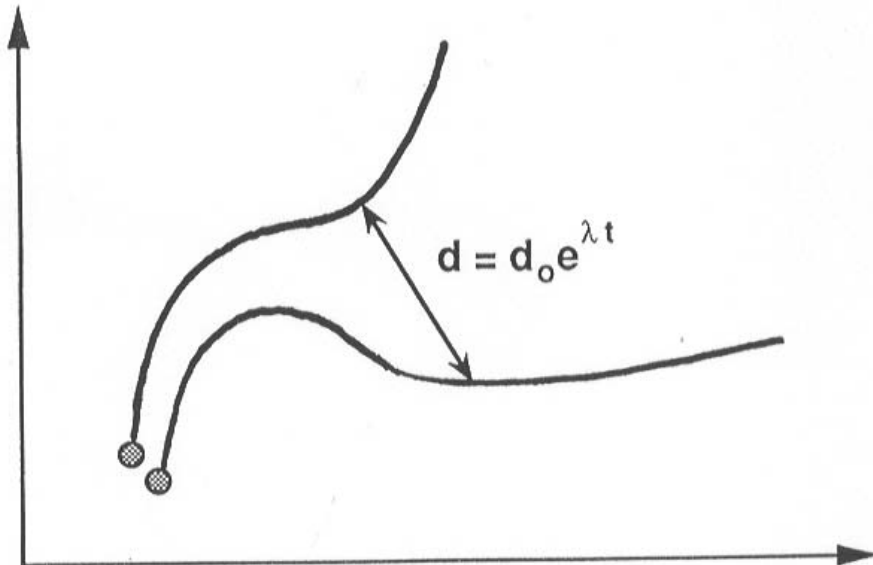
45. Au moyen du développement de R qui vient d'être donné, on pourra déterminer les valeurs de I, G, H, l, g, h, en fonction du temps, en se servant des équations (9). Les valeurs de ces six quantités devront ensuite être substituées dans les expressions des coordonnées de la Lune, ce qui donnera définitivement ces coordonnées en fonction du temps.

- **Neptune** was discovered by **Leverrier** (1811-1877) and **Adams** (1819-1892) using perturbation theory, due to anomalies observed in the motion of Uranus.
- What happens to the long-term stability of the planets?

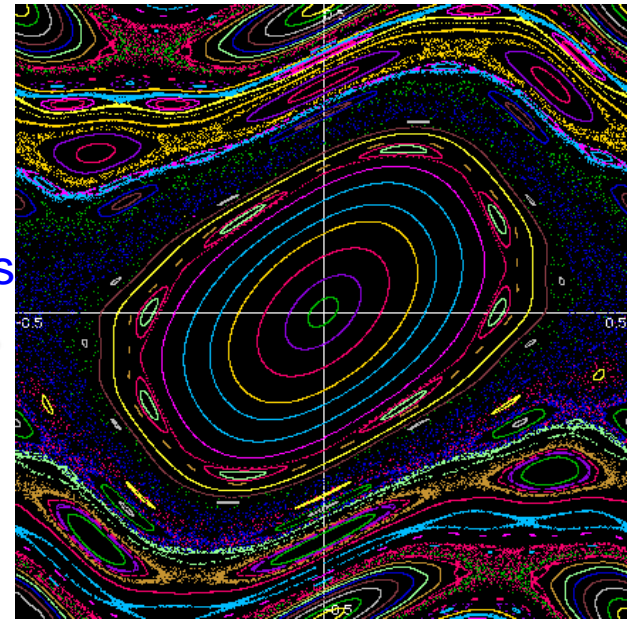
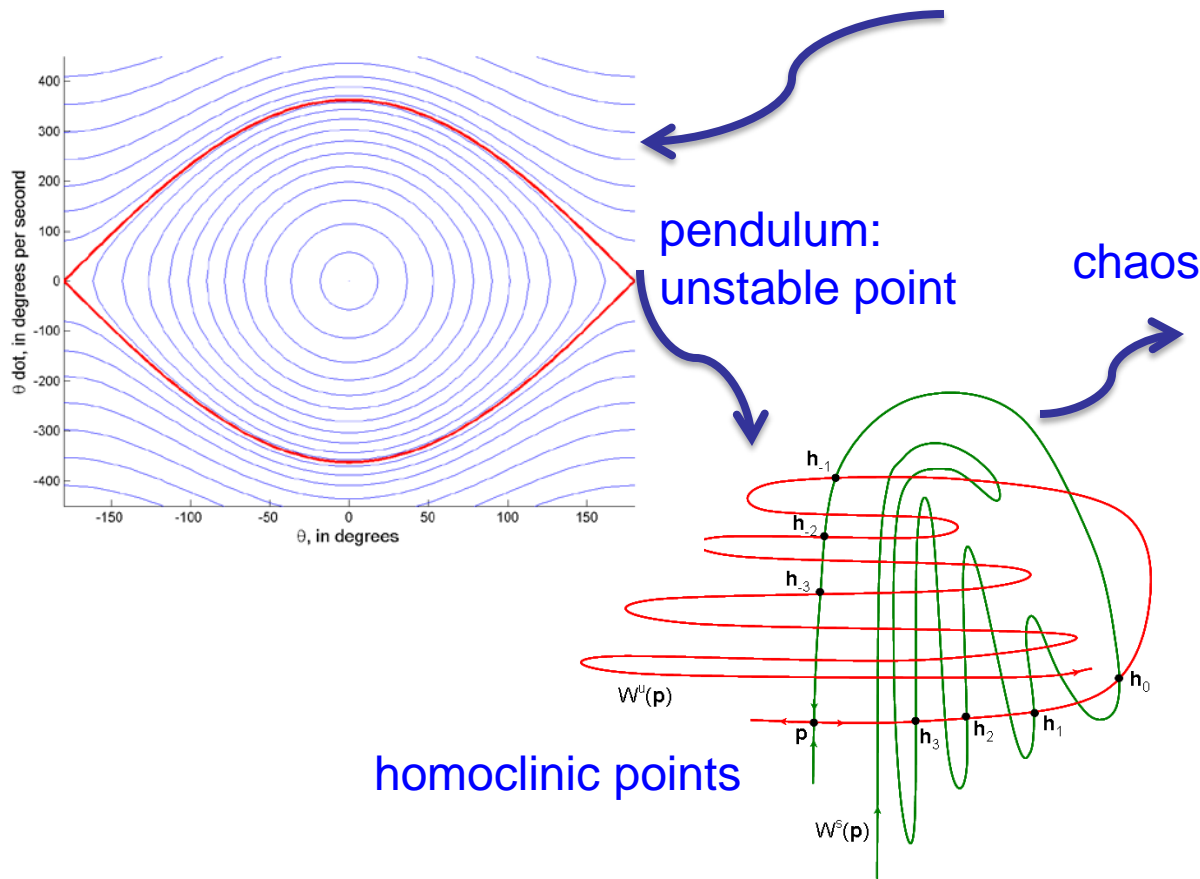


7. Chaos

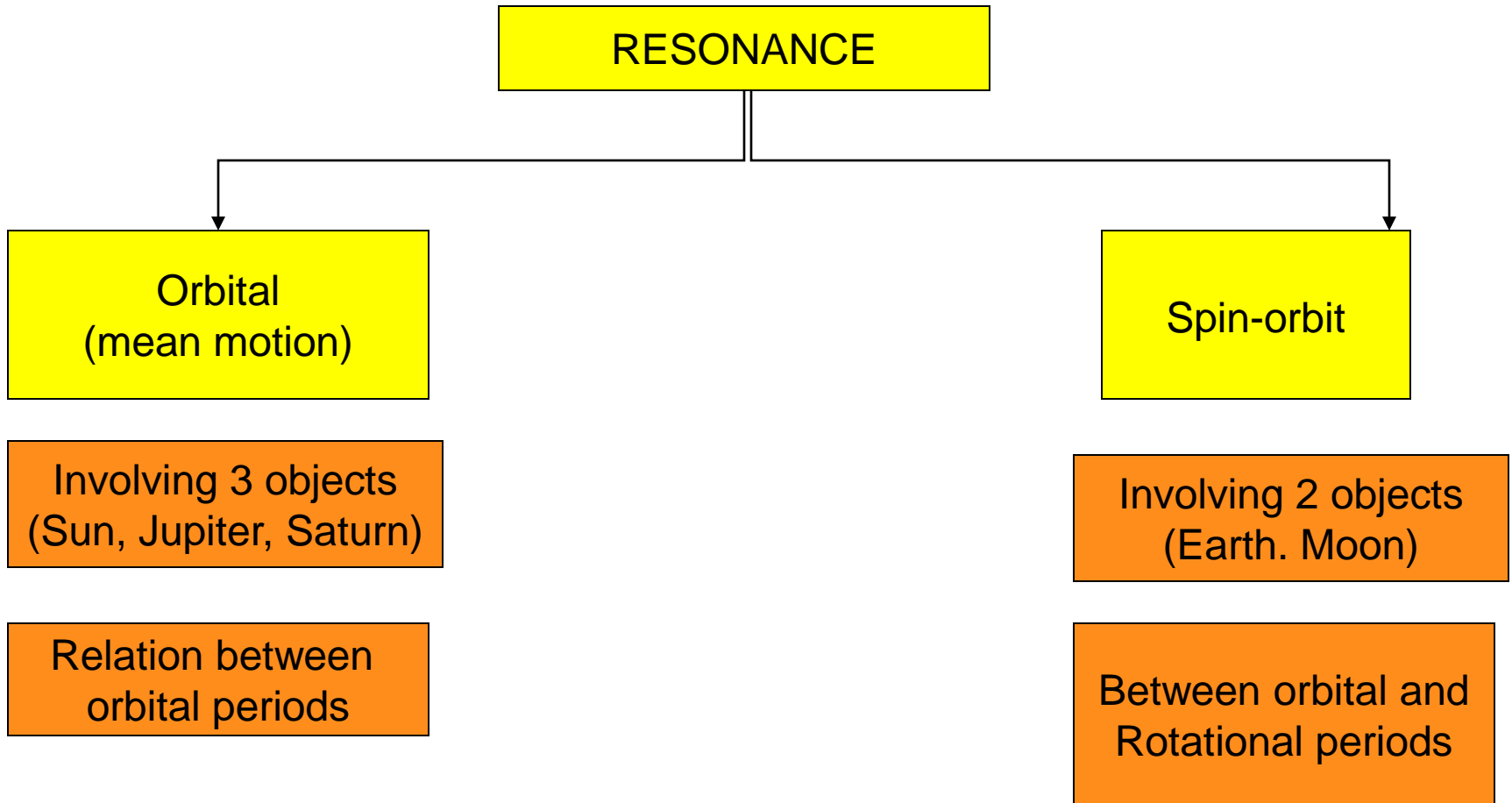
- Chaos: irregular motion showing an *extreme sensitivity to the choice of the initial conditions*.
- **Poincaré**: discovered chaos while studying the 3-body problem (later Lorenz in 1962 the “Butterfly Effect”).
- *Chaos does not mean that a system is unstable, but rather unpredictable.*



- Earth-Moon-spacecraft= 3-body pb, no Kepler laws
- Poincaré: 3-body problem, homoclinic points, chaos
- Kolmogorov: KAM theory, regular orbits



8. Resonances

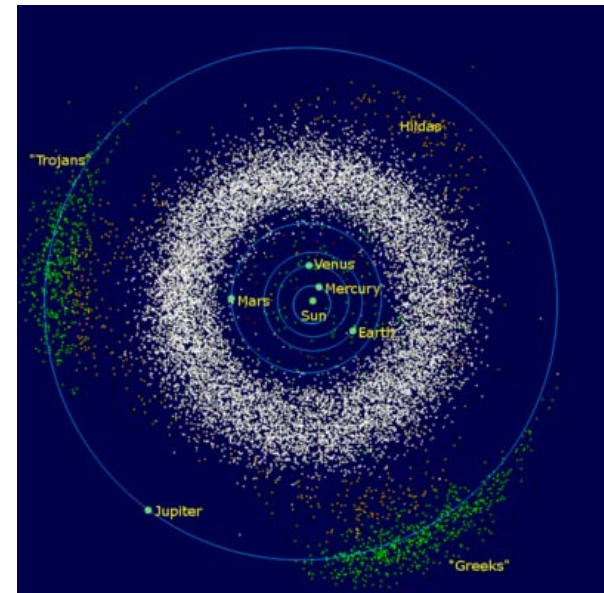
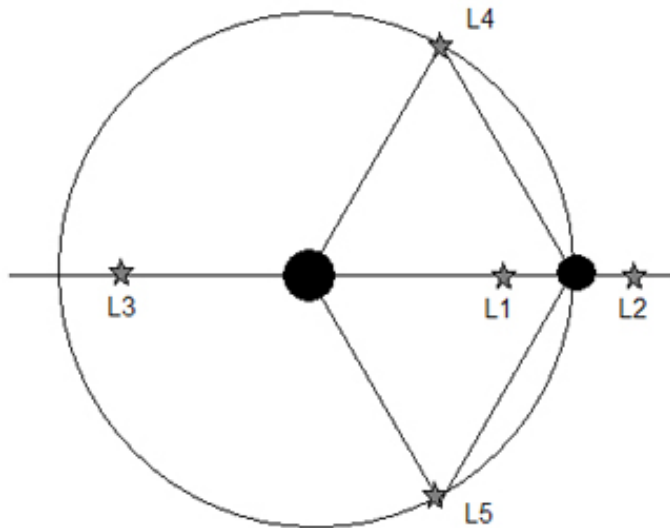


9. Orbital resonances

- 3 bodies: S (Sun), A (asteroid), J (Jupiter)
- Let T_A e T_J be the periods of revolution around S.
- **Definition:** An *orbital resonance* between A e J occurs when:
$$T_A / T_J = p/q \quad \text{with } p, q \text{ non-zero intergers.}$$
- **Examples:**
- **Jupiter and Saturn:** $T_J / T_S = 2/5$ or 2 Saturn's orbits correspond to 5 Jupiter's orbits;
- **Io, Europa, Ganimede, Callisto:** $T_{IO}/T_{EUR} = 1/2$,
 $T_{IO}/T_{GAN} = 1/4$, $T_{EUR}/T_{GAN} = 1/2$;
- **Satellites of Saturn:** $T_{Titan}/T_{Hyperion} = 3/4$, $T_{Titan}/T_{Japetus} = 1/5$;
- **Greek and Trojan asteroids** 1/1.

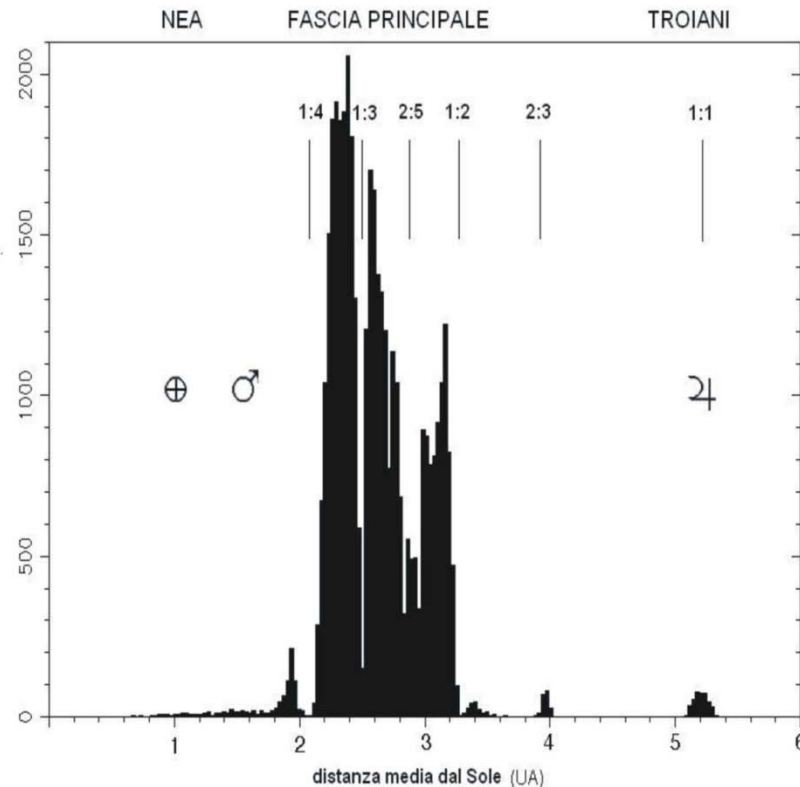
10. Greek and Trojan asteroids

- Two groups of asteroids in 1:1 resonance with Jupiter (same orbital period, same distance from the Sun).
- Euler collinear points L1, L2, L3; Lagrange triangular points L4, L5 (Greek and Trojans).



11. Full and empty resonances

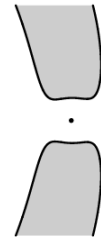
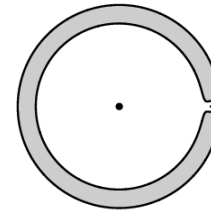
- Main belt asteroids between Mars and Jupiter: some resonances are full (1:1, 2:3), other regions called *Kirkwood gaps* are empty (1:2, 1:3, 1:4).



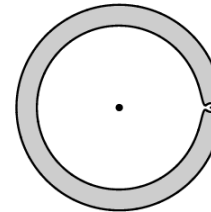
12. Interplanetary highways

- J.-L. Lagrange: Cette recherche n'est à la vérité que de pure curiosité
- C. Conley (1968): use the bottleneck between the primaries and chaos around the collinear points to travel at low cost (use Moser's version of Lyapunov theorem).

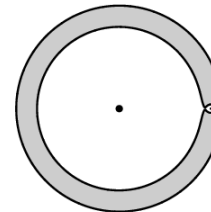
$$C_J = 3.03$$



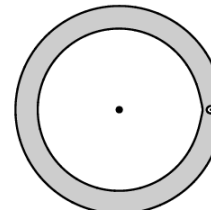
$$C_J = 3.038$$



$$C_J = 3.0395$$

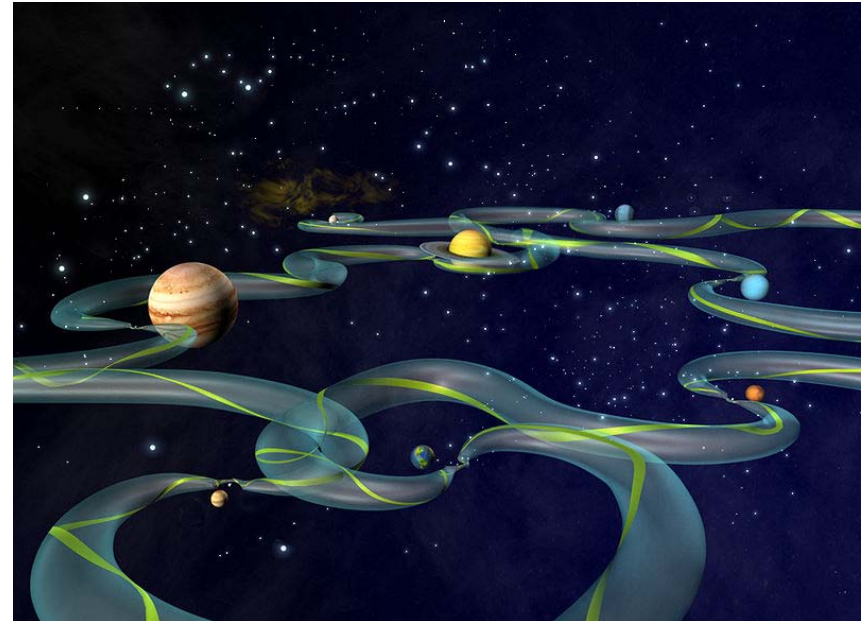


$$C_J = 3.045$$



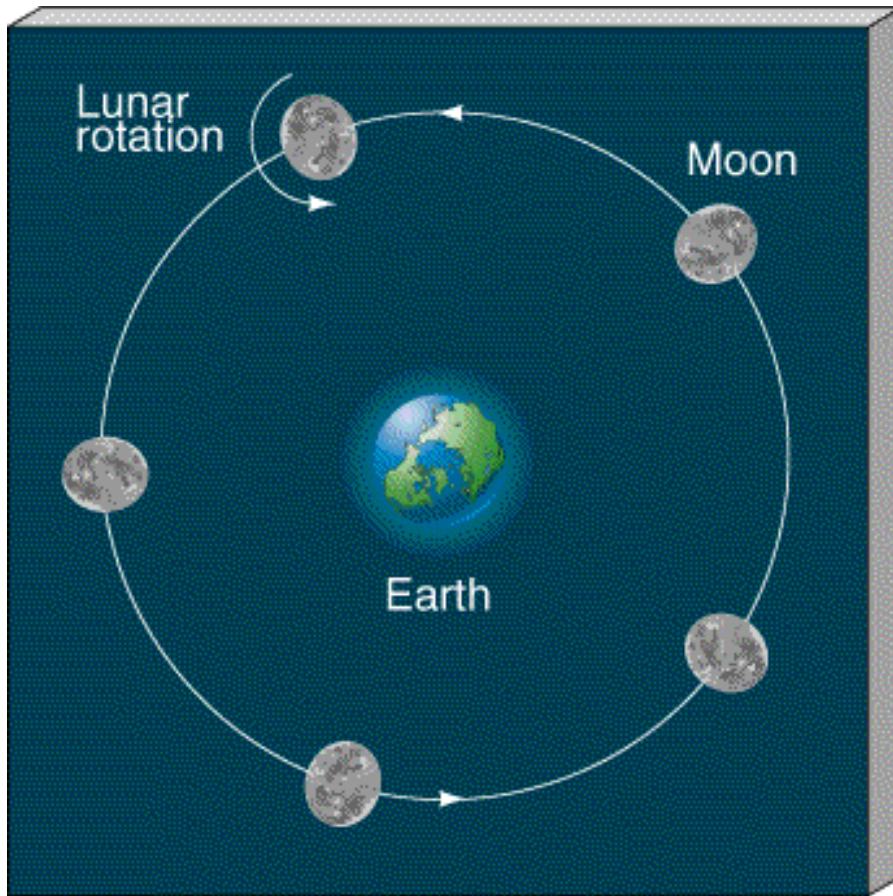
Conley: *“Unfortunately, orbits such as these require a long time to complete a cycle (e.g., 6 months, though a modification of the notion might improve that). On the other hand, one cannot predict how knowledge will be applied – only that it often is”.*

- International Sun/Earth Explorer 3 (ISEE-3) 1978
- SOHO (1995)
- MAP (2001)
- GENESIS (2001)
- HERSCHEL-PLANCK (2009)

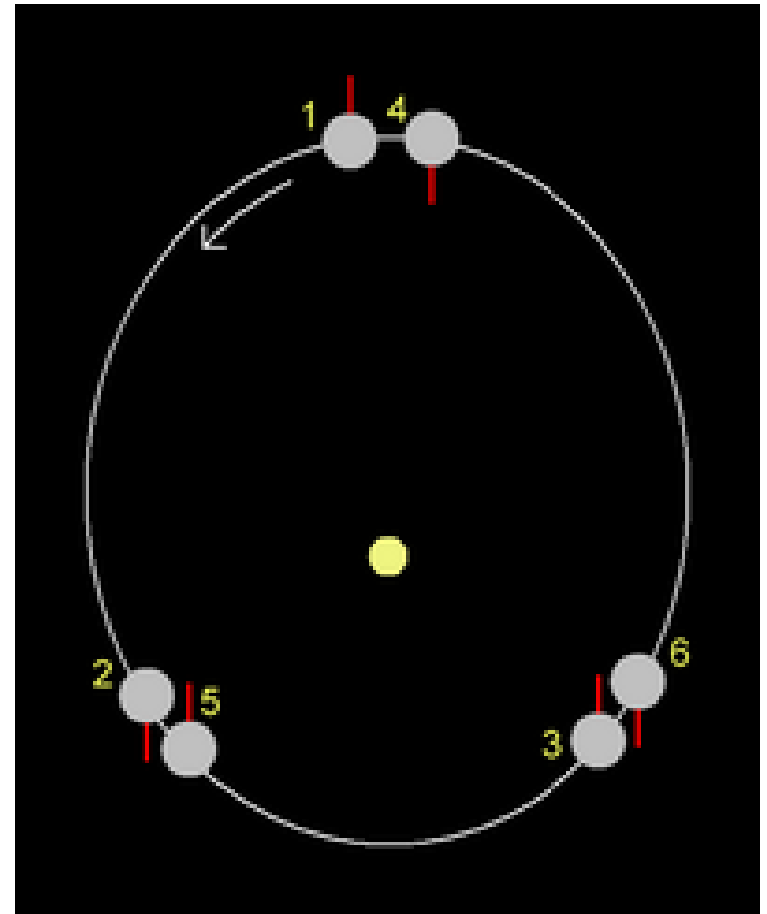


13. Spin-orbit resonances

- Earth and Moon with masses m_E e m_M ; T_{rev} orbital period of M around E and T_{rot} rotational period of M (rigid body) around an internal spin axis.
- **Definition**: A *spin-orbit resonance* of order p/q , occurs if
$$T_{\text{rev}} / T_{\text{rot}} = p/q$$
with p, q non-zero integers.
- Most famous example: 1/1 Earth-Moon spin-orbit resonance, where the Moon always points the same face to the Earth.
- Mercury-Sun: 3/2 spin-orbit resonance, 2 revolutions of Mercury around the Sun correspond to 3 rotations about its spin-axis.
- Hyperion is in chaotic spin-orbit dynamics.

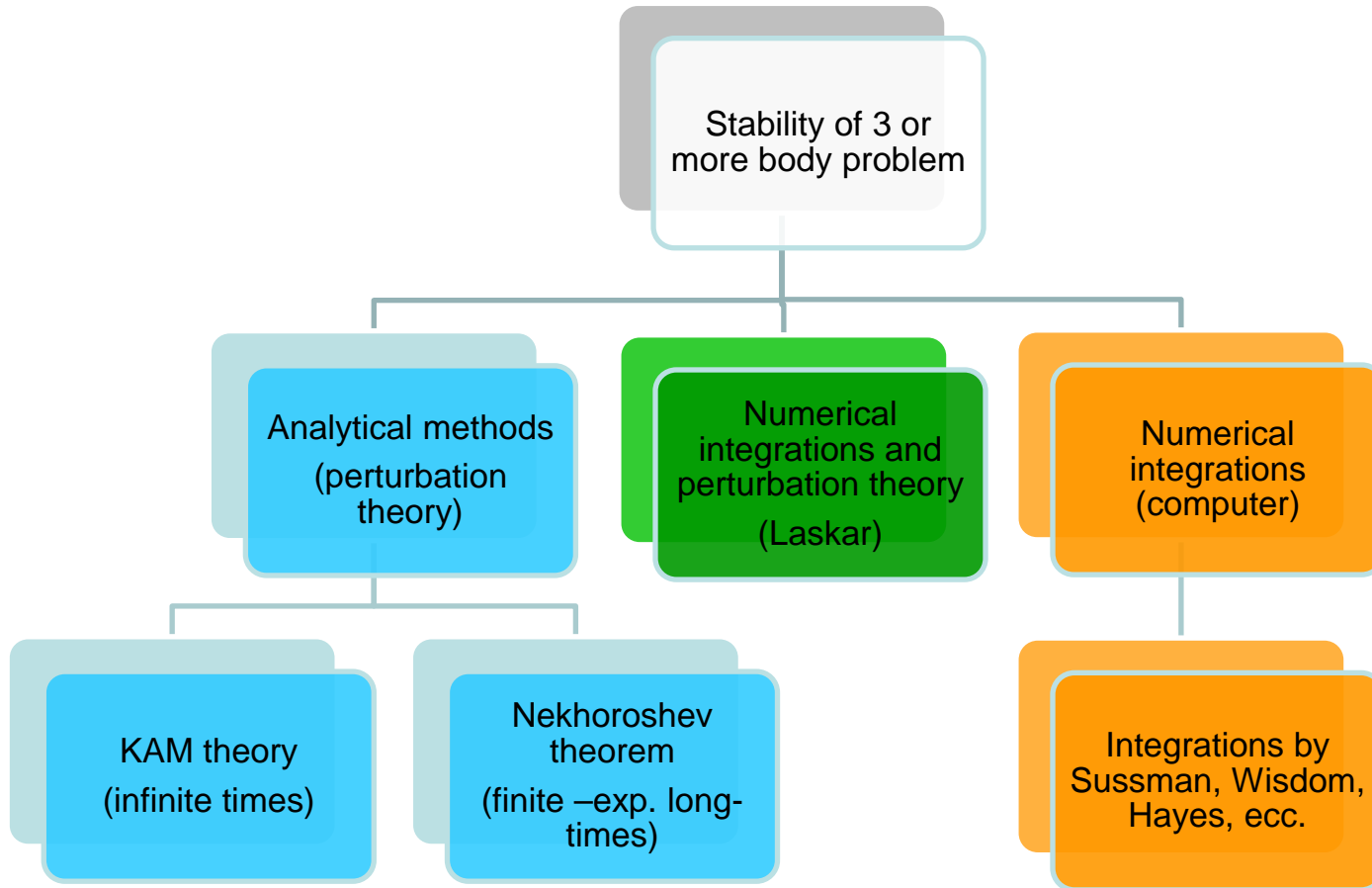


Resonance 1/1



Resonance 3/2

14. Is the Solar system stable?



- Laskar: the internal Solar system is **CHAOTIC**.
- From an error of **15 mt** on the initial position of the Earth: error **150 mt** after 10 million years
- error **150 million km** after 100 million years, **no further predictions!**

□ **RESULTS:**

- **Mercury** and **Mars** very chaotic
- **Venus** and **Earth** moderately chaotic
- **External planets** are regular
- **Pluto** very chaotic.

